

An Introduction to Three-Dimensional, Rigid Body Dynamics

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Volume I: Kinematics

Unit 7

Application of Concepts to Systems with Complex Interconnecting Joints

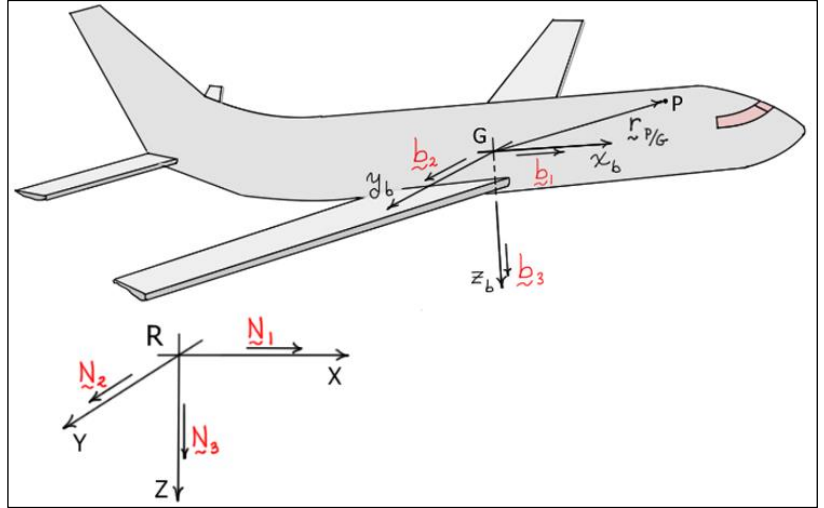
Summary

The *basic kinematical concepts* for mechanical systems were presented in Units 1 through 4. The concepts were applied to systems in which the *relative orientation changes* between adjoining bodies were described by a *single rotation* - simple revolute joints. In Units 5 and 6 concepts were presented for describing the relative orientation changes between *any two adjoining bodies* by using a set of up to three *orientation angles* or a set of four *Euler parameters*. This unit *applies* the concepts from Units 1 through 6 to analyze the kinematics of systems with *more complex connecting joints*. The major differences in applying the kinematical concepts of Units 1 through 4 are in the calculations of orientation, angular velocity and angular acceleration.

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Example 1:

The orientation of an aircraft relative to reference frame R is defined by a 3-2-1 **body-fixed** rotation sequence (ψ, θ, ϕ) . The **velocity** of the aircraft is given in body-fixed directions as ${}^R \underline{v}_G = u \underline{b}_1 + v \underline{b}_2 + w \underline{b}_3$. The **position** of a point P on the aircraft relative to the mass center G is given as $\underline{r}_{P/G} = p_1 \underline{b}_1 + p_2 \underline{b}_2 + p_3 \underline{b}_3$. Point P is assumed to be **fixed** on the aircraft.



Find:

- the components of ${}^R \underline{v}_P$ the **velocity** of P in R in both the **body-fixed** and **base** frames
- the components of ${}^R \underline{a}_P$ the **acceleration** of P in R in both the **body-fixed** and **base** frames

Solution:

- Using the method of **direct differentiation** as presented in Unit 2, ${}^R \underline{v}_P$ can be found as follows

$${}^R \underline{v}_P = \frac{{}^R d}{dt} (\underline{r}_P) = \frac{{}^R d}{dt} (\underline{r}_G + \underline{r}_{P/G}) = \frac{{}^R d}{dt} (\underline{r}_G) + \frac{{}^R d}{dt} (\underline{r}_{P/G}) = {}^R \underline{v}_G + \underbrace{\frac{{}^B d}{dt} (\underline{r}_{P/G})}_{\text{zero}} + ({}^R \underline{\omega}_B \times \underline{r}_{P/G})$$

$$= (u \underline{b}_1 + v \underline{b}_2 + w \underline{b}_3) + \begin{vmatrix} \underline{b}_1 & \underline{b}_2 & \underline{b}_3 \\ \omega_1 & \omega_2 & \omega_3 \\ p_1 & p_2 & p_3 \end{vmatrix}$$

$$\Rightarrow {}^R \underline{v}_P = \underbrace{(u + \omega_2 p_3 - \omega_3 p_2)}_{v_1} \underline{b}_1 + \underbrace{(v + \omega_3 p_1 - \omega_1 p_3)}_{v_2} \underline{b}_2 + \underbrace{(w + \omega_1 p_2 - \omega_2 p_1)}_{v_3} \underline{b}_3$$

Here (as presented in Unit 5): $\omega_1 = \dot{\phi} - \dot{\psi} S_\theta$, $\omega_2 = \dot{\theta} C_\phi + \dot{\psi} C_\theta S_\phi$, $\omega_3 = -\dot{\theta} S_\phi + \dot{\psi} C_\theta C_\phi$.

The same results can also be found by using the concepts for points **fixed on bodies** as presented in Unit 3. Using the **body-fixed** components of the **angular velocity** of the aircraft developed in Unit 5, the velocity of P may be written as

$${}^R \underline{v}_P = {}^R \underline{v}_G + {}^R \underline{v}_{P/G} = {}^R \underline{v}_G + \underbrace{({}^R \underline{\omega}_B \times \underline{r}_{P/G})}_{{}^R \underline{v}_{P/G}} = (u \underline{b}_1 + v \underline{b}_2 + w \underline{b}_3) + \begin{vmatrix} \underline{b}_1 & \underline{b}_2 & \underline{b}_3 \\ \omega_1 & \omega_2 & \omega_3 \\ p_1 & p_2 & p_3 \end{vmatrix}$$

This clearly produces the **same result** as above.

The components (V_i) of ${}^R \underline{v}_P$ in the **base frame** R may be found from the **body-fixed** components v_i ($i=1,2,3$) (as defined above) using the coordinate transformation matrix for a 3-2-1 orientation angle sequence derived in Unit 5.

$$\begin{Bmatrix} V_1 \\ V_2 \\ V_3 \end{Bmatrix} = [R]^T \begin{Bmatrix} v_1 \\ v_2 \\ v_3 \end{Bmatrix} = \begin{bmatrix} C_\psi C_\theta & S_\psi C_\theta & -S_\theta \\ C_\psi S_\theta S_\phi - S_\psi C_\phi & S_\psi S_\theta S_\phi + C_\psi C_\phi & C_\theta S_\phi \\ C_\psi S_\theta C_\phi + S_\psi S_\phi & S_\psi S_\theta C_\phi - C_\psi S_\phi & C_\theta C_\phi \end{bmatrix}^T \begin{Bmatrix} v_1 \\ v_2 \\ v_3 \end{Bmatrix}$$

b) Using the method of **direct differentiation** (along with the derivative rule) as presented in Unit 2, ${}^R \underline{a}_P$ can be found as follows

$$\begin{aligned} {}^R \underline{a}_P &= \frac{{}^R d}{dt} ({}^R \underline{v}_P) = \frac{{}^R d}{dt} ({}^R \underline{v}_P) = \frac{{}^B d}{dt} ({}^R \underline{v}_P) + ({}^R \underline{\omega}_B \times {}^R \underline{v}_P) \\ &= (\dot{v}_1 \underline{b}_1 + \dot{v}_2 \underline{b}_2 + \dot{v}_3 \underline{b}_3) + \begin{vmatrix} \underline{b}_1 & \underline{b}_2 & \underline{b}_3 \\ \omega_1 & \omega_2 & \omega_3 \\ v_1 & v_2 & v_3 \end{vmatrix} \\ &= (\dot{u} + \dot{\omega}_2 p_3 - \dot{\omega}_3 p_2 + \omega_2 v_3 - \omega_3 v_2) \underline{b}_1 + (\dot{v} + \dot{\omega}_3 p_1 - \dot{\omega}_1 p_3 + \omega_3 v_1 - \omega_1 v_3) \underline{b}_2 \\ &\quad + (\dot{w} + \dot{\omega}_1 p_2 - \dot{\omega}_2 p_1 + \omega_1 v_2 - \omega_2 v_1) \underline{b}_3 \\ &= (\dot{u} + \alpha_2 p_3 - \alpha_3 p_2 + \omega_2 (w + \omega_1 p_2 - \omega_2 p_1) - \omega_3 (v + \omega_3 p_1 - \omega_1 p_3)) \underline{b}_1 \\ &\quad + (\dot{v} + \alpha_3 p_1 - \alpha_1 p_3 + \omega_3 (u + \omega_2 p_3 - \omega_3 p_2) - \omega_1 (w + \omega_1 p_2 - \omega_2 p_1)) \underline{b}_2 \\ &\quad + (\dot{w} + \alpha_1 p_2 - \alpha_2 p_1 + \omega_1 (v + \omega_3 p_1 - \omega_1 p_3) - \omega_2 (u + \omega_2 p_3 - \omega_3 p_2)) \underline{b}_3 \end{aligned}$$

Or

$$\boxed{{}^R \underline{a}_P = a_1 \underline{b}_1 + a_2 \underline{b}_2 + a_3 \underline{b}_3}$$

where

$$\begin{aligned} a_1 &= \dot{u} + \alpha_2 p_3 - \alpha_3 p_2 + \omega_2 (w + \omega_1 p_2 - \omega_2 p_1) - \omega_3 (v + \omega_3 p_1 - \omega_1 p_3) \\ a_2 &= \dot{v} + \alpha_3 p_1 - \alpha_1 p_3 + \omega_3 (u + \omega_2 p_3 - \omega_3 p_2) - \omega_1 (w + \omega_1 p_2 - \omega_2 p_1) \\ a_3 &= \dot{w} + \alpha_1 p_2 - \alpha_2 p_1 + \omega_1 (v + \omega_3 p_1 - \omega_1 p_3) - \omega_2 (u + \omega_2 p_3 - \omega_3 p_2) \end{aligned}$$

This result can also be found using the concepts for **points fixed** on bodies as presented in Unit 3. Using the **body-fixed** components of the **angular velocity** and **angular acceleration** vectors of the aircraft developed in Unit 5, the **acceleration** of P may be written as

$${}^R \underline{a}_P = {}^R \underline{a}_G + {}^R \underline{a}_{P/G} = \frac{{}^R d}{dt} ({}^R \underline{v}_G) + ({}^R \underline{\alpha}_B \times \underline{r}_{P/G}) + ({}^R \underline{\omega}_B \times {}^R \underline{v}_{P/G})$$

Using the *derivative rule* for the first term gives

$$\begin{aligned}
 {}^R \underline{a}_P &= \frac{B}{dt} \left({}^R \underline{v}_G \right) + \left({}^R \underline{\omega}_B \times {}^R \underline{v}_G \right) + \left({}^R \underline{\alpha}_B \times \underline{r}_{P/G} \right) + \left({}^R \underline{\omega}_B \times {}^R \underline{v}_{P/G} \right) \\
 &= \left(\dot{u} \underline{b}_1 + \dot{v} \underline{b}_2 + \dot{w} \underline{b}_3 \right) + \begin{vmatrix} \underline{b}_1 & \underline{b}_2 & \underline{b}_3 \\ \omega_1 & \omega_2 & \omega_3 \\ u & v & w \end{vmatrix} + \begin{vmatrix} \underline{b}_1 & \underline{b}_2 & \underline{b}_3 \\ \alpha_1 & \alpha_2 & \alpha_3 \\ p_1 & p_2 & p_3 \end{vmatrix} \\
 &\quad + \begin{vmatrix} \underline{b}_1 & \underline{b}_2 & \underline{b}_3 \\ \omega_1 & \omega_2 & \omega_3 \\ (\omega_2 p_3 - \omega_3 p_2) & (\omega_3 p_1 - \omega_1 p_3) & (\omega_1 p_2 - \omega_2 p_1) \end{vmatrix} \\
 &\Rightarrow \boxed{{}^R \underline{a}_P = a_1 \underline{b}_1 + a_2 \underline{b}_2 + a_3 \underline{b}_3}
 \end{aligned}$$

Expanding the determinants and combining terms gives the same components a_i ($i=1,2,3$) found above using direct differentiation. The angular velocity components ω_i ($i=1,2,3$) are as defined in part (a), and the angular acceleration components α_i ($i=1,2,3$) are as computed in Unit 5.

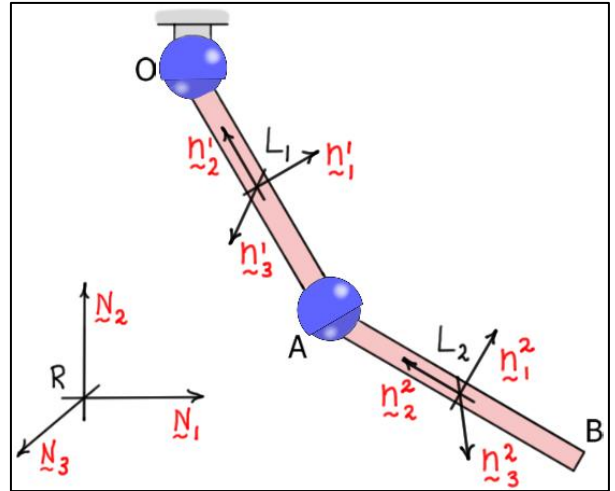
$$\begin{aligned}
 \alpha_1 &= \dot{\omega}_1 = \ddot{\phi} - \dot{\psi} S_\theta - \dot{\psi} \dot{\theta} C_\theta \\
 \alpha_2 &= \dot{\omega}_2 = \ddot{\theta} C_\phi - \dot{\theta} \dot{\phi} S_\phi + \dot{\psi} C_\theta S_\phi - \dot{\psi} \dot{\theta} S_\theta S_\phi + \dot{\psi} \dot{\phi} C_\theta C_\phi \\
 \alpha_3 &= \dot{\omega}_3 = -\ddot{\theta} S_\phi - \dot{\theta} \dot{\phi} C_\phi + \dot{\psi} C_\theta C_\phi - \dot{\psi} \dot{\theta} S_\theta C_\phi - \dot{\psi} \dot{\phi} C_\theta S_\phi
 \end{aligned}$$

The components (A_i) of ${}^R \underline{a}_P$ in the *base frame* R may be found from the *body-fixed* components (a_i as defined above) using the *coordinate transformation matrix*.

$$\begin{aligned}
 \begin{Bmatrix} A_1 \\ A_2 \\ A_3 \end{Bmatrix} &= [R]^T \begin{Bmatrix} a_1 \\ a_2 \\ a_3 \end{Bmatrix} = \begin{bmatrix} C_\psi C_\theta & S_\psi C_\theta & -S_\theta \\ C_\psi S_\theta S_\phi - S_\psi C_\phi & S_\psi S_\theta S_\phi + C_\psi C_\phi & C_\theta S_\phi \\ C_\psi S_\theta C_\phi + S_\psi S_\phi & S_\psi S_\theta C_\phi - C_\psi S_\phi & C_\theta C_\phi \end{bmatrix}^T \begin{Bmatrix} v_1 \\ v_2 \\ v_3 \end{Bmatrix}
 \end{aligned}$$

Example 2:

The system shown is a **three-dimensional double pendulum** or **arm**. The first link is connected to ground and the second link is connected to the first with ball and socket joints at O and A . The ground frame is $R: (\underline{N}_1, \underline{N}_2, \underline{N}_3)$ and the link frames are $L_i: (\underline{n}_1^i, \underline{n}_2^i, \underline{n}_3^i)$ ($i=1,2$). The **orientation** of each link is defined relative to R using a 3-1-3 **body-fixed** rotation sequence. The **lengths** of the links are ℓ_1 and ℓ_2 .



Find:

- the components of ${}^R \underline{v}_B$ the **velocity** of B in R in the **base** frame
- the components of ${}^R \underline{a}_B$ the **acceleration** of B in R in the **base** frame

Solution:

- Using the method of **direct differentiation** (along with the derivative rule) as presented in Unit 2, the velocity ${}^R \underline{v}_B$ can be found as follows

$$\begin{aligned}
 {}^R \underline{v}_B &= \frac{{}^R d}{dt} (\underline{r}_B) = \frac{{}^R d}{dt} (\underline{r}_{A/O} + \underline{r}_{B/A}) = \frac{{}^R d}{dt} (-\ell_1 \underline{n}_2^1) + \frac{{}^R d}{dt} (-\ell_2 \underline{n}_2^2) \\
 &= \underbrace{\frac{{}^{L_1} d}{dt} (-\ell_1 \underline{n}_2^1)}_{\text{zero}} + \left({}^R \underline{\omega}_{L_1} \times -\ell_1 \underline{n}_2^1 \right) + \underbrace{\frac{{}^{L_2} d}{dt} (-\ell_2 \underline{n}_2^2)}_{\text{zero}} + \left({}^R \underline{\omega}_{L_2} \times -\ell_2 \underline{n}_2^2 \right) \\
 &= \begin{vmatrix} \underline{n}_1^1 & \underline{n}_2^1 & \underline{n}_3^1 \\ \omega_{11} & \omega_{12} & \omega_{13} \\ 0 & -\ell_1 & 0 \end{vmatrix} + \begin{vmatrix} \underline{n}_1^2 & \underline{n}_2^2 & \underline{n}_3^2 \\ \omega_{21} & \omega_{22} & \omega_{23} \\ 0 & -\ell_2 & 0 \end{vmatrix} \\
 &= \left((\ell_1 \omega_{13}) \underline{n}_1^1 + (0) \underline{n}_2^1 + (-\ell_1 \omega_{11}) \underline{n}_3^1 \right) + \left((\ell_2 \omega_{23}) \underline{n}_1^2 + (0) \underline{n}_2^2 + (-\ell_2 \omega_{21}) \underline{n}_3^2 \right) \\
 &= \begin{bmatrix} \ell_1 \omega_{13} & 0 & -\ell_1 \omega_{11} \end{bmatrix} \begin{Bmatrix} \underline{n}_1^1 \\ \underline{n}_2^1 \\ \underline{n}_3^1 \end{Bmatrix} + \begin{bmatrix} \ell_2 \omega_{23} & 0 & -\ell_2 \omega_{21} \end{bmatrix} \begin{Bmatrix} \underline{n}_1^2 \\ \underline{n}_2^2 \\ \underline{n}_3^2 \end{Bmatrix} \\
 &= \begin{bmatrix} \ell_1 \omega_{13} & 0 & -\ell_1 \omega_{11} \end{bmatrix} [R_1] \begin{Bmatrix} \underline{N}_1 \\ \underline{N}_2 \\ \underline{N}_3 \end{Bmatrix} + \begin{bmatrix} \ell_2 \omega_{23} & 0 & -\ell_2 \omega_{21} \end{bmatrix} [R_2] \begin{Bmatrix} \underline{N}_1 \\ \underline{N}_2 \\ \underline{N}_3 \end{Bmatrix}
 \end{aligned}$$

So, the components of ${}^R\mathbf{v}_B$ in the base system are

$$\begin{aligned} \begin{Bmatrix} V_1 \\ V_2 \\ V_3 \end{Bmatrix} &= \begin{bmatrix} V_1 & V_2 & V_3 \end{bmatrix}^T = \begin{bmatrix} \ell_1\omega_{13} & 0 & -\ell_1\omega_{11} \end{bmatrix} \begin{bmatrix} R_1 \end{bmatrix} + \begin{bmatrix} \ell_2\omega_{23} & 0 & -\ell_2\omega_{21} \end{bmatrix} \begin{bmatrix} R_2 \end{bmatrix}^T \\ &= \begin{bmatrix} \ell_1\omega_{13} & 0 & -\ell_1\omega_{11} \end{bmatrix} \begin{bmatrix} R_1 \end{bmatrix}^T + \begin{bmatrix} \ell_2\omega_{23} & 0 & -\ell_2\omega_{21} \end{bmatrix} \begin{bmatrix} R_2 \end{bmatrix}^T \\ \Rightarrow \begin{Bmatrix} V_1 \\ V_2 \\ V_3 \end{Bmatrix} &= \begin{bmatrix} R_1 \end{bmatrix}^T \begin{Bmatrix} \ell_1\omega_{13} \\ 0 \\ -\ell_1\omega_{11} \end{Bmatrix} + \begin{bmatrix} R_2 \end{bmatrix}^T \begin{Bmatrix} \ell_2\omega_{23} \\ 0 \\ -\ell_2\omega_{21} \end{Bmatrix} \end{aligned}$$

The **transformation matrices** of the links and the **body-fixed components** of the **angular velocity vectors** are as given in Unit 5.

$$\begin{aligned} \boxed{\begin{bmatrix} R_i \end{bmatrix}} &= \begin{bmatrix} C_{i1}C_{i3} - S_{i1}C_{i2}S_{i3} & S_{i1}C_{i3} + C_{i1}C_{i2}S_{i3} & S_{i2}S_{i3} \\ -C_{i1}S_{i3} - S_{i1}C_{i2}C_{i3} & -S_{i1}S_{i3} + C_{i1}C_{i2}C_{i3} & S_{i2}C_{i3} \\ S_{i1}S_{i2} & -C_{i1}S_{i2} & C_{i2} \end{bmatrix} & \boxed{\begin{matrix} \omega_{i1} = \dot{\theta}_{i1}S_{i2}S_{i3} + \dot{\theta}_{i2}C_{i3} \\ \omega_{i2} = \dot{\theta}_{i1}S_{i2}C_{i3} - \dot{\theta}_{i2}S_{i3} \\ \omega_{i3} = \dot{\theta}_{i3} + \dot{\theta}_{i1}C_{i2} \end{matrix}} \quad (i=1,2) \end{aligned}$$

The same results can also be found by using the concepts for points **fixed** on bodies as presented in Unit 3. Using the **body-fixed components** of the **angular velocities** of the links given in Unit 5, the **velocity** of B may be written as

$$\begin{aligned} {}^R\mathbf{v}_B &= {}^R\mathbf{v}_A + {}^R\mathbf{v}_{B/A} = {}^R\mathbf{v}_{A/O} + {}^R\mathbf{v}_{B/A} = \underbrace{\left({}^R\boldsymbol{\omega}_{L_1} \times \mathbf{r}_{A/O} \right)}_{{}^R\mathbf{v}_{A/O}} + \underbrace{\left({}^R\boldsymbol{\omega}_{L_2} \times \mathbf{r}_{B/A} \right)}_{{}^R\mathbf{v}_{B/A}} = \begin{vmatrix} \underline{n}_1 & \underline{n}_2 & \underline{n}_3 \\ \omega_{11} & \omega_{12} & \omega_{13} \\ 0 & -\ell_1 & 0 \end{vmatrix} + \begin{vmatrix} \underline{n}_1 & \underline{n}_2 & \underline{n}_3 \\ \omega_{21} & \omega_{22} & \omega_{23} \\ 0 & -\ell_2 & 0 \end{vmatrix} \\ &= \dots \end{aligned}$$

This is clearly going to produce the **same results** as above.

b) Using the method of **direct differentiation** (along with the derivative rule) as presented in Unit 2, the acceleration ${}^R\mathbf{a}_P$ can be found as follows

$$\begin{aligned} {}^R\mathbf{a}_B &= \frac{{}^R d}{{}^R dt} \left({}^R\mathbf{v}_B \right) = \frac{{}^R d}{{}^R dt} \left((\ell_1\omega_{13})\underline{n}_1^1 + (0)\underline{n}_2^1 + (-\ell_1\omega_{11})\underline{n}_3^1 \right) + \frac{{}^R d}{{}^R dt} \left((\ell_2\omega_{23})\underline{n}_1^2 + (0)\underline{n}_2^2 + (-\ell_2\omega_{21})\underline{n}_3^2 \right) \\ &= \frac{{}^R d}{{}^R dt} \left((\ell_1\omega_{13})\underline{n}_1^1 + (0)\underline{n}_2^1 + (-\ell_1\omega_{11})\underline{n}_3^1 \right) + \left({}^R\boldsymbol{\omega}_{L_1} \times \left((\ell_1\omega_{13})\underline{n}_1^1 + (0)\underline{n}_2^1 + (-\ell_1\omega_{11})\underline{n}_3^1 \right) \right) \\ &\quad + \frac{{}^R d}{{}^R dt} \left((\ell_2\omega_{23})\underline{n}_1^2 + (0)\underline{n}_2^2 + (-\ell_2\omega_{21})\underline{n}_3^2 \right) + \left({}^R\boldsymbol{\omega}_{L_2} \times \left((\ell_2\omega_{23})\underline{n}_1^2 + (0)\underline{n}_2^2 + (-\ell_2\omega_{21})\underline{n}_3^2 \right) \right) \end{aligned}$$

Expanding the cross products using determinants gives

$$\begin{aligned}
{}^R \underline{a}_B &= \left((\ell_1 \dot{\omega}_{13}) \underline{n}_1^1 + (0) \underline{n}_2^1 + (-\ell_1 \dot{\omega}_{11}) \underline{n}_3^1 \right) + \left((\ell_2 \dot{\omega}_{23}) \underline{n}_1^2 + (0) \underline{n}_2^2 + (-\ell_2 \dot{\omega}_{21}) \underline{n}_3^2 \right) \\
&+ \begin{vmatrix} \underline{n}_1^1 & \underline{n}_2^1 & \underline{n}_3^1 \\ \omega_{11} & \omega_{12} & \omega_{13} \\ \ell_1 \omega_{13} & 0 & -\ell_1 \omega_{11} \end{vmatrix} + \begin{vmatrix} \underline{n}_1^2 & \underline{n}_2^2 & \underline{n}_3^2 \\ \omega_{21} & \omega_{22} & \omega_{23} \\ \ell_2 \omega_{23} & 0 & -\ell_2 \omega_{21} \end{vmatrix} \\
&= \left(\ell_1 \alpha_{13} - \ell_1 \omega_{11} \omega_{12} \right) \underline{n}_1^1 + \left(\ell_1 \omega_{13}^2 + \ell_1 \omega_{11}^2 \right) \underline{n}_2^1 + \left(-\ell_1 \alpha_{11} - \ell_1 \omega_{12} \omega_{13} \right) \underline{n}_3^1 \\
&+ \left(\ell_2 \alpha_{23} - \ell_2 \omega_{21} \omega_{22} \right) \underline{n}_1^2 + \left(\ell_2 \omega_{23}^2 + \ell_2 \omega_{21}^2 \right) \underline{n}_2^2 + \left(-\ell_2 \alpha_{21} - \ell_2 \omega_{22} \omega_{23} \right) \underline{n}_3^2
\end{aligned}$$

Part of this result is expressed in the L_1 frame and part is expressed in the L_2 frame. Using the transformation matrices for the links, the components in the base frame may be written as

$$\begin{bmatrix} A_1 \\ A_2 \\ A_3 \end{bmatrix} = [R_1]^T \begin{bmatrix} \ell_1 \alpha_{13} - \ell_1 \omega_{11} \omega_{12} \\ \ell_1 \omega_{13}^2 + \ell_1 \omega_{11}^2 \\ -\ell_1 \alpha_{11} - \ell_1 \omega_{12} \omega_{13} \end{bmatrix} + [R_2]^T \begin{bmatrix} \ell_2 \alpha_{23} - \ell_2 \omega_{21} \omega_{22} \\ \ell_2 \omega_{23}^2 + \ell_2 \omega_{21}^2 \\ -\ell_2 \alpha_{21} - \ell_2 \omega_{22} \omega_{23} \end{bmatrix}$$

The **transformation matrices** and the **body-fixed components** of the **angular velocity vectors** are as defined in part (a), and the **body-fixed components** of the **angular acceleration vectors** are as given in Unit 5.

$$\begin{cases} \alpha_{i1} = \dot{\omega}_{i1} = \ddot{\theta}_{i1} S_{i2} S_{i3} + \dot{\theta}_{i1} \dot{\theta}_{i2} C_{i2} S_{i3} + \dot{\theta}_{i1} \dot{\theta}_{i3} S_{i2} C_{i3} + \ddot{\theta}_{i2} C_{i3} - \dot{\theta}_{i2} \dot{\theta}_{i3} S_{i3} \\ \alpha_{i2} = \dot{\omega}_{i2} = \ddot{\theta}_{i1} S_{i2} C_{i3} + \dot{\theta}_{i1} \dot{\theta}_{i2} C_{i2} C_{i3} - \dot{\theta}_{i1} \dot{\theta}_{i3} S_{i2} S_{i3} - \ddot{\theta}_{i2} S_{i3} - \dot{\theta}_{i2} \dot{\theta}_{i3} C_{i3} \\ \alpha_{i3} = \dot{\omega}_{i3} = \ddot{\theta}_{i3} + \ddot{\theta}_{i1} C_{i2} - \dot{\theta}_{i1} \dot{\theta}_{i2} S_{i2} \end{cases} \quad (i=1,2)$$

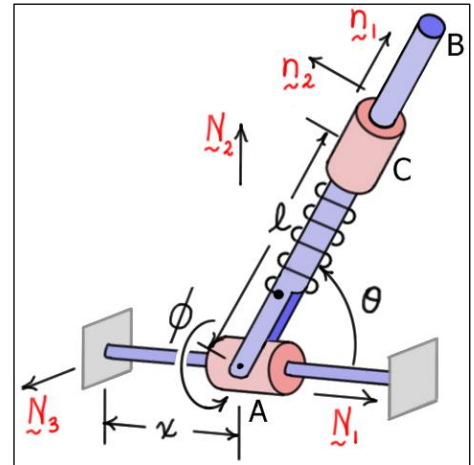
The same results can also be found by using the concepts for points **fixed on bodies** as presented in Unit 3. Using the **body-fixed** components of the **angular velocities** and **angular accelerations** of the links given in Unit 5, the acceleration of B may be written as

$$\begin{aligned}
{}^R \underline{a}_B &= {}^R \underline{a}_A + {}^R \underline{a}_{B/A} = {}^R \underline{a}_{A/O} + {}^R \underline{a}_{B/A} \\
&= \left({}^R \underline{\alpha}_{L_1} \times {}^R \underline{r}_{A/O} \right) + \left({}^R \underline{\omega}_{L_1} \times {}^R \underline{v}_{A/O} \right) + \left({}^R \underline{\alpha}_{L_2} \times {}^R \underline{r}_{B/A} \right) + \left({}^R \underline{\omega}_{L_2} \times {}^R \underline{v}_{B/A} \right) \\
&= \begin{vmatrix} \underline{n}_1^1 & \underline{n}_2^1 & \underline{n}_3^1 \\ \alpha_{11} & \alpha_{12} & \alpha_{13} \\ 0 & -\ell_1 & 0 \end{vmatrix} + \begin{vmatrix} \underline{n}_1^1 & \underline{n}_2^1 & \underline{n}_3^1 \\ \omega_{11} & \omega_{12} & \omega_{13} \\ \ell_1 \omega_{13} & 0 & -\ell_1 \omega_{11} \end{vmatrix} + \begin{vmatrix} \underline{n}_1^2 & \underline{n}_2^2 & \underline{n}_3^2 \\ \alpha_{21} & \alpha_{22} & \alpha_{23} \\ 0 & -\ell_2 & 0 \end{vmatrix} + \begin{vmatrix} \underline{n}_1^2 & \underline{n}_2^2 & \underline{n}_3^2 \\ \omega_{21} & \omega_{22} & \omega_{23} \\ \ell_2 \omega_{23} & 0 & -\ell_2 \omega_{21} \end{vmatrix} \\
&= \left(\ell_1 \alpha_{13} - \ell_1 \omega_{11} \omega_{12} \right) \underline{n}_1^1 + \left(\ell_1 \omega_{13}^2 + \ell_1 \omega_{11}^2 \right) \underline{n}_2^1 + \left(-\ell_1 \alpha_{11} - \ell_1 \omega_{12} \omega_{13} \right) \underline{n}_3^1 \\
&+ \left(\ell_2 \alpha_{23} - \ell_2 \omega_{21} \omega_{22} \right) \underline{n}_1^2 + \left(\ell_2 \omega_{23}^2 + \ell_2 \omega_{21}^2 \right) \underline{n}_2^2 + \left(-\ell_2 \alpha_{21} - \ell_2 \omega_{22} \omega_{23} \right) \underline{n}_3^2 \\
&= \dots
\end{aligned}$$

These are clearly the **same results** as found above.

Example 3:

The system shown consists of three bodies – collar A, bar AB, and collar C. Collar A *slides along* and *rotates about* the fixed horizontal bar. Its displacement (x) and rotation (ϕ) are both in the *fixed* \underline{N}_1 direction. This type of connection is often referred to as a *cylindrical joint*. Bar AB is pinned to collar A. Its *orientation* relative to the collar (and the horizontal bar) is defined by the angle θ . Finally, collar C is assumed to simply *translate* along the bar AB and not rotate about it. Its position relative to A is given by the length ℓ , and its motion along the bar is restricted by the attached spring.



Given this setup, the orientation of the bar frame $AB: (\underline{n}_1, \underline{n}_2, \underline{n}_3)$ relative to the ground frame $R: (\underline{N}_1, \underline{N}_2, \underline{N}_3)$ may be described by a 1-3 *body-fixed* rotation sequence. To orient the bar relative to the ground, first align the two frames. In this position, the bar lies along the \underline{N}_1 direction, and the pin axis at A is aligned with the \underline{N}_3 direction. Then rotate the bar first through an angle ϕ about the \underline{N}_1 direction (simulating the rotation of collar A), and then rotate the bar through an angle θ about the \underline{n}_3 direction (simulating the pin rotation at A).

Find:

- the components of ${}^R \underline{v}_C$ the *velocity* of C in R in the *ground* frame
- the components of ${}^R \underline{a}_C$ the *acceleration* of C in R in the *ground* frame

Solution:

a) Using the method of *direct differentiation* (along with the derivative rule) as presented in Unit 2, the velocity ${}^R \underline{v}_C$ can be found as follows

$$\begin{aligned} {}^R \underline{v}_C &= \frac{{}^R d}{dt}(\underline{r}_C) = \frac{{}^R d}{dt}(\underline{r}_A + \underline{r}_{C/A}) = \frac{{}^R d}{dt}(x \underline{N}_1 + \ell \underline{n}_1) = \frac{{}^R d}{dt}(x \underline{N}_1) + \frac{{}^R d}{dt}(\ell \underline{n}_1) \\ &= \dot{x} \underline{N}_1 + \frac{{}^{AB} d}{dt}(\ell \underline{n}_1) + ({}^R \underline{\omega}_{AB} \times \ell \underline{n}_1) = \dot{x} \underline{N}_1 + \dot{\ell} \underline{n}_1 + \begin{vmatrix} \underline{n}_1 & \underline{n}_2 & \underline{n}_3 \\ \omega_1 & \omega_2 & \omega_3 \\ \ell & 0 & 0 \end{vmatrix} \\ &= \dot{x} \underline{N}_1 + \dot{\ell} \underline{n}_1 + (\ell \omega_3) \underline{n}_2 + (-\ell \omega_2) \underline{n}_3 \end{aligned}$$

Using the transformation matrix for a 1-3 rotation sequence, the ground-frame components of ${}^R \underline{v}_C$ may be written as

$$\begin{Bmatrix} V_1 \\ V_2 \\ V_3 \end{Bmatrix} = \begin{Bmatrix} \dot{x} \\ 0 \\ 0 \end{Bmatrix} + [R]^T \begin{Bmatrix} \dot{l} \\ \ell \omega_3 \\ -\ell \omega_2 \end{Bmatrix}$$

The **transformation matrix** and the **body-fixed components** of the **angular velocity** vector may be found using the techniques presented in Unit 5. First, align the unit vectors $(\underline{n}_1, \underline{n}_2, \underline{n}_3)$ with the ground-frame unit vectors $(\underline{N}_1, \underline{N}_2, \underline{N}_3)$. Then, rotate about an angle ϕ about the $\underline{N}_1 = \underline{N}'_1$ direction, and then rotate about a second angle θ about the rotated $\underline{N}'_3 = \underline{n}_3$ direction. Using this procedure, the three sets of unit vectors for the orientations of bar AB may be related as follows.

$$\begin{Bmatrix} \underline{N}_1 \\ \underline{N}_2 \\ \underline{N}_3 \end{Bmatrix} = [R_1]^T \begin{Bmatrix} \underline{N}'_1 \\ \underline{N}'_2 \\ \underline{N}'_3 \end{Bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & C_\phi & S_\phi \\ 0 & -S_\phi & C_\phi \end{bmatrix}^T \begin{Bmatrix} \underline{N}'_1 \\ \underline{N}'_2 \\ \underline{N}'_3 \end{Bmatrix} \quad \begin{Bmatrix} \underline{N}'_1 \\ \underline{N}'_2 \\ \underline{N}'_3 \end{Bmatrix} = [R_2]^T \begin{Bmatrix} \underline{n}_1 \\ \underline{n}_2 \\ \underline{n}_3 \end{Bmatrix} = \begin{bmatrix} C_\theta & S_\theta & 0 \\ -S_\theta & C_\theta & 0 \\ 0 & 0 & 1 \end{bmatrix}^T \begin{Bmatrix} \underline{n}_1 \\ \underline{n}_2 \\ \underline{n}_3 \end{Bmatrix}$$

Recall from Unit 5 that C_ϕ and C_θ represent the cosines of angles ϕ and θ , and S_ϕ and S_θ represent the sines of angles ϕ and θ . The transformation matrix for bar AB is then found to be

$$[R] = [R_2][R_1] = \begin{bmatrix} C_\theta & S_\theta & 0 \\ -S_\theta & C_\theta & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 \\ 0 & C_\phi & S_\phi \\ 0 & -S_\phi & C_\phi \end{bmatrix} = \begin{bmatrix} C_\theta & C_\phi S_\theta & S_\phi S_\theta \\ -S_\theta & C_\phi C_\theta & S_\phi C_\theta \\ 0 & -S_\phi & C_\phi \end{bmatrix}$$

As used above, the matrix $[R]^T$ transforms components of vector expressed in AB into components in the ground frame.

As in Unit 5, the **body-fixed** components of the **angular velocity vectors** of the links are found using the **summation rule** for angular velocities.

$${}^R\omega_{AB} = \dot{\phi} \underline{N}'_1 + \dot{\theta} \underline{n}_3 = \dot{\phi} (C_\theta \underline{n}_1 - S_\theta \underline{n}_2) + \dot{\theta} \underline{n}_3 \Rightarrow \begin{cases} \omega_1 = \dot{\phi} C_\theta \\ \omega_2 = -\dot{\phi} S_\theta \\ \omega_3 = \dot{\theta} \end{cases}$$

The same results can be found using the concepts for **points moving on bodies** as presented in Unit 4. Using the **body-fixed components** of the **angular velocity** of bar AB , the **velocity** of C may be written as

$${}^R\mathcal{V}_C = {}^R\mathcal{V}_{\hat{C}} + {}^{AB}\mathcal{V}_C = {}^R\mathcal{V}_A + {}^R\mathcal{V}_{\hat{C}/A} + {}^{AB}\mathcal{V}_C \quad (C \text{ moves on } AB, \hat{C} \text{ is fixed on } AB \text{ and coincides with } C)$$

where

$$\begin{aligned} {}^R \underline{v}_{\hat{C}} &= {}^R \underline{v}_A + {}^R \underline{v}_{\hat{C}/A} = \dot{x} N_1 + \left({}^R \underline{\omega}_{AB} \times \underline{r}_{\hat{C}/A} \right) = \dot{x} N_2 + \left(\omega_1 \underline{n}_1 + \omega_2 \underline{n}_2 + \omega_3 \underline{n}_3 \right) \times \left(\ell \underline{n}_1 \right) \\ &= \dot{x} N_1 + \ell \left(\omega_3 \underline{n}_2 - \omega_2 \underline{n}_3 \right) \end{aligned}$$

$${}^{AB} \underline{v}_C = \dot{\ell} \underline{n}_1 \quad (C \text{ has } \mathbf{rectilinear motion} \text{ on } AB)$$

Substituting into the boxed equation gives

$${}^R \underline{v}_C = \dot{x} N_1 + \dot{\ell} \underline{n}_1 + \left(\ell \omega_3 \right) \underline{n}_2 + \left(-\ell \omega_2 \right) \underline{n}_3 \quad \dots \text{ same results as above}$$

b) Using the method of **direct differentiation** (along with the derivative rule) as presented in Unit 2, the acceleration ${}^R \underline{a}_C$ can be found as follows

$$\begin{aligned} {}^R \underline{a}_C &= \frac{{}^R d}{dt} \left(\dot{x} N_1 + \dot{\ell} \underline{n}_1 + \ell \omega_3 \underline{n}_2 - \ell \omega_2 \underline{n}_3 \right) = \frac{{}^R d}{dt} \left(\dot{x} N_1 \right) + \frac{{}^R d}{dt} \left(\dot{\ell} \underline{n}_1 + \ell \omega_3 \underline{n}_2 - \ell \omega_2 \underline{n}_3 \right) \\ &= \ddot{x} N_1 + \frac{{}^{AB} d}{dt} \left(\dot{\ell} \underline{n}_1 + \ell \omega_3 \underline{n}_2 - \ell \omega_2 \underline{n}_3 \right) + {}^R \underline{\omega}_{AB} \times \left(\dot{\ell} \underline{n}_1 + \ell \omega_3 \underline{n}_2 - \ell \omega_2 \underline{n}_3 \right) \\ &= \ddot{x} N_1 + \left(\ddot{\ell} \underline{n}_1 + \left(\dot{\ell} \omega_3 + \ell \dot{\omega}_3 \right) \underline{n}_2 - \left(\dot{\ell} \omega_2 + \ell \dot{\omega}_2 \right) \underline{n}_3 \right) + \begin{vmatrix} \underline{n}_1 & \underline{n}_2 & \underline{n}_3 \\ \omega_1 & \omega_2 & \omega_3 \\ \dot{\ell} & \ell \omega_3 & -\ell \omega_2 \end{vmatrix} \\ &= \ddot{x} N_1 + \left(\ddot{\ell} - \ell \omega_2^2 - \ell \omega_3^2 \right) \underline{n}_1 + \left(\dot{\ell} \omega_3 + \ell \dot{\omega}_3 + \dot{\ell} \omega_3 + \ell \omega_1 \omega_2 \right) \underline{n}_2 + \left(-\dot{\ell} \omega_2 - \ell \dot{\omega}_2 + \ell \omega_1 \omega_3 - \dot{\ell} \omega_2 \right) \underline{n}_3 \\ &= \ddot{x} N_1 + \left(\ddot{\ell} - \ell \omega_2^2 - \ell \omega_3^2 \right) \underline{n}_1 + \left(2\dot{\ell} \omega_3 + \ell \dot{\omega}_3 + \ell \omega_1 \omega_2 \right) \underline{n}_2 + \left(-2\dot{\ell} \omega_2 - \ell \dot{\omega}_2 + \ell \omega_1 \omega_3 \right) \underline{n}_3 \end{aligned}$$

Using the transformation matrix defined above, the ground-frame components of ${}^R \underline{a}_C$ are

$$\begin{bmatrix} A_1 \\ A_2 \\ A_3 \end{bmatrix} = \begin{bmatrix} \ddot{x} \\ 0 \\ 0 \end{bmatrix} + [R]^T \begin{bmatrix} \ddot{\ell} - \ell \left(\omega_2^2 + \omega_3^2 \right) \\ 2\dot{\ell} \omega_3 + \ell \left(\alpha_3 + \omega_1 \omega_2 \right) \\ \ell \left(\omega_1 \omega_3 - \alpha_2 \right) - 2\dot{\ell} \omega_2 \end{bmatrix} \quad \text{with} \quad \begin{bmatrix} \alpha_1 = \dot{\omega}_1 = \ddot{\phi} C_\theta - \dot{\phi} \dot{\theta} S_\theta \\ \alpha_2 = \dot{\omega}_2 = -\ddot{\phi} S_\theta - \dot{\phi} \dot{\theta} C_\theta \\ \alpha_3 = \dot{\omega}_3 = \ddot{\theta} \end{bmatrix}$$

The same results may be found by using the concepts for **points moving on bodies** as presented in Unit 4. Using the **body-fixed components** of the **angular velocity** and **angular acceleration** of bar AB, the **acceleration** of C may be written as

$$\boxed{{}^R \underline{a}_C = {}^R \underline{a}_{\hat{C}} + {}^{AB} \underline{a}_C + 2 \left({}^R \underline{\omega}_{AB} \times {}^{AB} \underline{v}_C \right)} \quad (C \text{ moves on } AB, \hat{C} \text{ is fixed on } AB \text{ and coincides with } C)$$

where, using the **two-point formula**,

$${}^R \underline{a}_{\hat{C}} = {}^R \underline{a}_A + {}^R \underline{a}_{\hat{C}/A} = \ddot{x} N_1 + \left({}^R \underline{\omega}_{AB} \times \underline{r}_{C/A} \right) + \left({}^R \underline{\omega}_{AB} \times {}^R \underline{v}_{\hat{C}/A} \right)$$

Expanding each of the terms gives

$$\begin{aligned}
{}^R \underline{a}_C &= \ddot{x} \underline{N}_1 + \left[(\alpha_1 \underline{n}_1 + \alpha_2 \underline{n}_2 + \alpha_3 \underline{n}_3) \times (\ell \underline{n}_1) \right] + \begin{vmatrix} \underline{n}_1 & \underline{n}_2 & \underline{n}_3 \\ \omega_1 & \omega_2 & \omega_3 \\ 0 & \ell \omega_3 & -\ell \omega_2 \end{vmatrix} \\
&= \ddot{x} \underline{N}_1 + \left[\ell (\alpha_3 \underline{n}_2 - \alpha_2 \underline{n}_3) \right] + \left[-\ell (\omega_2^2 + \omega_3^2) \underline{n}_1 + \ell \omega_1 \omega_2 \underline{n}_2 + \ell \omega_1 \omega_3 \underline{n}_3 \right] \\
&= \ddot{x} \underline{N}_1 + \left[-\ell (\omega_2^2 + \omega_3^2) \underline{n}_1 + \ell (\alpha_3 + \omega_1 \omega_2) \underline{n}_2 + \ell (\omega_1 \omega_3 - \alpha_2) \underline{n}_3 \right]
\end{aligned}$$

$${}^{AB} \underline{a}_C = \ddot{\ell} \underline{n}_1 \quad (C \text{ has } \textit{rectilinear motion} \text{ on } AB)$$

$$2 \left({}^R \underline{\omega}_{AB} \times {}^{AB} \underline{v}_C \right) = 2 \left(\omega_1 \underline{n}_1 + \omega_2 \underline{n}_2 + \omega_3 \underline{n}_3 \right) \times (\dot{\ell} \underline{n}_1) = 2 \dot{\ell} (\omega_3 \underline{n}_2 - \omega_2 \underline{n}_3) \quad (\textit{Coriolis acceleration})$$

Substituting these results into the boxed equation gives

$$\boxed{{}^R \underline{a}_C = \ddot{x} \underline{N}_1 + \left[\left(\ddot{\ell} - \ell (\omega_2^2 + \omega_3^2) \right) \underline{n}_1 + \left(2 \dot{\ell} \omega_3 + \ell (\alpha_3 + \omega_1 \omega_2) \right) \underline{n}_2 + \left(\ell (\omega_1 \omega_3 - \alpha_2) - 2 \dot{\ell} \omega_2 \right) \underline{n}_3 \right]}$$

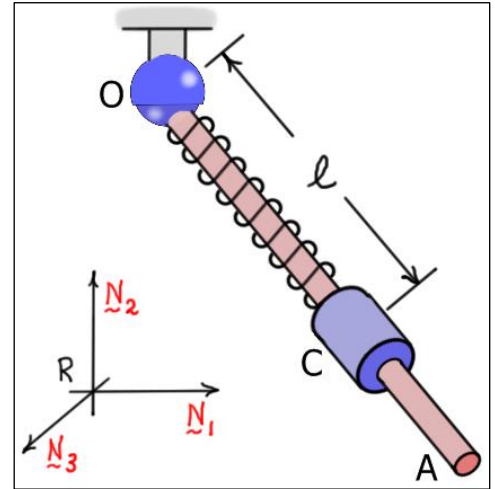
These are the *same results* as found above.

Notes:

1. In Example 2, the orientation of bar *AB* is described *relative to the ground-frame* rather than to the *adjoining bar OA*. In this case, this is easy to do because the interconnecting joint is spherical (ball-in-socket) – the *analytical model* of the connection is not limited by the choice of intermediate reference frames. If, however, the joint is a *two-axis universal joint*, then the analytical model must *accurately depict* the directions about which rotation is allowed. In this case, one of these directions is *fixed relative to the first body*, and the other is fixed *relative to the second*. In cases like this, it may be more meaningful (or convenient) to describe the orientation of a body relative to the adjoining body rather than the ground. See Exercise 7.2.
2. If orientation angles are measured *relative* to a *ground* (inertial) frame, they are referred to as *absolute* angles. If they are measured *relative* to an *adjoining body*, they are referred to as *relative* angles.
3. In Example 3, collar *A* is connected to the ground using a joint that allows *translation* and *rotation* of the collar *about a single axis*. This type of joint is referred to as a *cylindrical* joint. As the joint connects bar *AB* to the ground, the two coordinates (x, ϕ) are *absolute* coordinates. If, however, this type of joint is used to connect two bodies within a system, it may be more *natural* to define the coordinates as measured *relative* to the adjoining body. In this way, they will be more conveniently used to describe the *relative* motions of the two bodies.
4. In the above examples and in the Exercises below, the transformation matrices and angular velocity components are written in terms of *orientation angles* and their derivatives. In each case, however, they could have been written in terms of a set of four *Euler parameters* and their derivatives. See Unit 6.

Exercises:

7.1 The system shown consists of a bar OA and a collar C . Bar OA is connected to the ground using a ball-in-socket joint at O . As the bar swings, collar C translates along the bar. The variable length ℓ defines the position of C relative to O . The orientation of the bar relative to the ground-frame $R: (\underline{N}_1, \underline{N}_2, \underline{N}_3)$ is described by 3-1-3 body-fixed rotation sequence. When the orientation angles are all zero, the bar hangs vertically downward. Find a) the ground frame components of ${}^R \underline{v}_C$ the velocity of collar C relative to the ground frame, and b) the ground-frame components of ${}^R \underline{a}_C$ the acceleration of collar C relative to the ground frame.

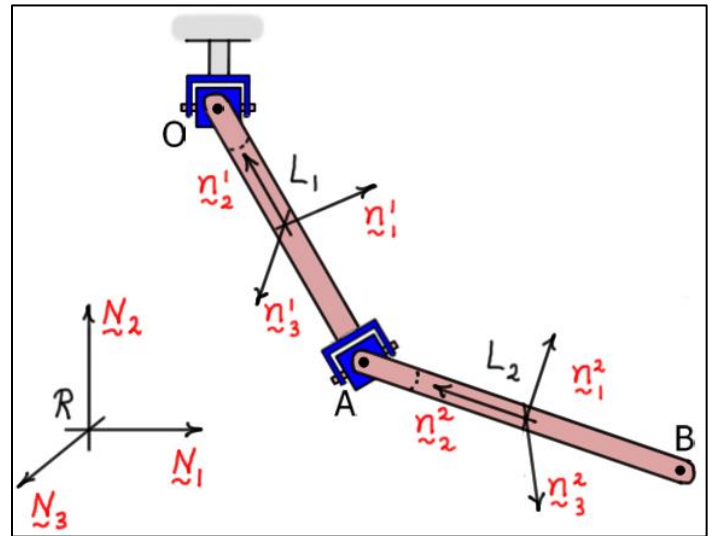


Answers: (using **body-fixed** angular velocity components)

$$\begin{Bmatrix} V_1 \\ V_2 \\ V_3 \end{Bmatrix} = [R]^T \begin{Bmatrix} \ell \omega_3 \\ -\dot{\ell} \\ -\ell \omega_1 \end{Bmatrix} \quad [R] = \begin{bmatrix} C_1 C_3 - S_1 C_2 S_3 & S_1 C_3 + C_1 C_2 S_3 & S_2 S_3 \\ -C_1 S_3 - S_1 C_2 C_3 & -S_1 S_3 + C_1 C_2 C_3 & S_2 C_3 \\ S_1 S_2 & -C_1 S_2 & C_2 \end{bmatrix} \quad \begin{cases} \omega_1 = \dot{\theta}_1 S_2 S_3 + \dot{\theta}_2 C_3 \\ \omega_2 = \dot{\theta}_1 S_2 C_3 - \dot{\theta}_2 S_3 \\ \omega_3 = \dot{\theta}_3 + \dot{\theta}_1 C_2 \end{cases}$$

$$\begin{Bmatrix} A_1 \\ A_2 \\ A_3 \end{Bmatrix} = [R]^T \begin{Bmatrix} 2\dot{\ell}\omega_3 + \ell(\alpha_3 - \omega_1\omega_2) \\ -\ddot{\ell} + \ell(\omega_1^2 + \omega_2^2) \\ -2\dot{\ell}\omega_1 - \ell(\alpha_1 + \omega_2\omega_3) \end{Bmatrix} \quad \begin{cases} \alpha_1 = \dot{\omega}_1 = \ddot{\theta}_1 S_2 S_3 + \dot{\theta}_1 \dot{\theta}_2 C_2 S_3 + \dot{\theta}_1 \dot{\theta}_3 S_2 C_3 + \ddot{\theta}_2 C_3 - \dot{\theta}_2 \dot{\theta}_3 S_3 \\ \alpha_2 = \dot{\omega}_2 = \ddot{\theta}_1 S_2 C_3 + \dot{\theta}_1 \dot{\theta}_2 C_2 C_3 - \dot{\theta}_1 \dot{\theta}_3 S_2 S_3 - \ddot{\theta}_2 S_3 - \dot{\theta}_2 \dot{\theta}_3 C_3 \\ \alpha_3 = \dot{\omega}_3 = \ddot{\theta}_3 + \ddot{\theta}_1 C_2 - \dot{\theta}_1 \dot{\theta}_2 S_2 \end{cases}$$

7.2 The system shown is a **three-dimensional double pendulum** or **arm**. The first link is connected to ground and the second link is connected to the first with **universal joints** at O and A , respectively. The ground frame is $R: (\underline{N}_1, \underline{N}_2, \underline{N}_3)$ and the link frames are $L_i: (\underline{n}_1^i, \underline{n}_2^i, \underline{n}_3^i)$ ($i=1,2$). The **orientation** of L_1 is defined **relative** to R and the orientation of L_2 is defined **relative** to L_1 each with a 1-3 **body-fixed** rotation sequence.



Link OA is oriented relative to the ground frame by first rotating through an angle θ_{11} about the \underline{n}_1 direction, and then rotating about an angle θ_{12} about the \underline{n}_3^1 direction. Link AB is oriented relative to link OA by rotating first through an angle θ_{21} about the \underline{n}_1^1 direction, and then through an angle θ_{22} about the \underline{n}_3^2 direction. The **lengths** of the links are ℓ_1 and ℓ_2 . Find a) the components of ${}^R \underline{v}_B$ the **velocity** of B in R in the **ground** frame, and b) the components of ${}^R \underline{a}_B$ the **acceleration** of B in R in the **ground** frame.

Answers: (using **body-fixed** angular velocity components)

$$\begin{Bmatrix} V_1 \\ V_2 \\ V_3 \end{Bmatrix} = [R_1]^T \begin{Bmatrix} \ell_1 \omega_{13} \\ 0 \\ -\ell_1 \omega_{11} \end{Bmatrix} + [R_1]^T [R_2]^T \begin{Bmatrix} \ell_2 \omega_{23} \\ 0 \\ -\ell_2 \omega_{21} \end{Bmatrix}$$

$$\begin{Bmatrix} A_1 \\ A_2 \\ A_3 \end{Bmatrix} = [R_1]^T \begin{Bmatrix} \ell_1 \alpha_{13} - \ell_1 \omega_{11} \omega_{12} \\ \ell_1 \omega_{13}^2 + \ell_1 \omega_{11}^2 \\ -\ell_1 \alpha_{11} - \ell_1 \omega_{12} \omega_{13} \end{Bmatrix} + [R_1]^T [R_2]^T \begin{Bmatrix} \ell_2 \alpha_{23} - \ell_2 \omega_{21} \omega_{22} \\ \ell_2 \omega_{23}^2 + \ell_2 \omega_{21}^2 \\ -\ell_2 \alpha_{21} - \ell_2 \omega_{22} \omega_{23} \end{Bmatrix}$$

$$[R_i] = \begin{bmatrix} C_{i2} & S_{i2} & 0 \\ -S_{i2} & C_{i2} & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 \\ 0 & C_{i1} & S_{i1} \\ 0 & -S_{i1} & C_{i1} \end{bmatrix} = \begin{bmatrix} C_{i2} & C_{i1} S_{i2} & S_{i1} S_{i2} \\ -S_{i2} & C_{i1} C_{i2} & S_{i1} C_{i2} \\ 0 & -S_{i1} & C_{i1} \end{bmatrix} \quad (i=1,2) \dots \text{one for each link}$$

$${}^R \underline{\omega}_{L_1} = \omega_{11} \underline{n}_1^1 + \omega_{12} \underline{n}_2^1 + \omega_{13} \underline{n}_3^1 = \dot{\theta}_{11} (C_{12} \underline{n}_1^1 - S_{12} \underline{n}_2^1) + \dot{\theta}_{12} \underline{n}_3^1$$

$${}^R \underline{\omega}_{L_2} = {}^R \underline{\omega}_{L_1} + {}^{L_1} \underline{\omega}_{L_2} = \omega_{21} \underline{n}_1^2 + \omega_{22} \underline{n}_2^2 + \omega_{23} \underline{n}_3^2 \quad \begin{Bmatrix} \omega_{21} \\ \omega_{22} \\ \omega_{23} \end{Bmatrix} = [R_2] \begin{Bmatrix} C_{12} \dot{\theta}_{11} \\ -S_{12} \dot{\theta}_{11} \\ \dot{\theta}_{12} \end{Bmatrix} + \begin{Bmatrix} C_{22} \dot{\theta}_{21} \\ -S_{22} \dot{\theta}_{21} \\ \dot{\theta}_{22} \end{Bmatrix}$$

$$\begin{aligned} \alpha_{11} &= \dot{\omega}_{11} = C_{12} \ddot{\theta}_{11} - S_{12} \dot{\theta}_{11} \dot{\theta}_{12} \\ \alpha_{12} &= \dot{\omega}_{12} = -S_{12} \ddot{\theta}_{11} - C_{12} \dot{\theta}_{11} \dot{\theta}_{12} \\ \alpha_{13} &= \dot{\omega}_{13} = \ddot{\theta}_{12} \end{aligned}$$

$$\begin{Bmatrix} \alpha_{21} \\ \alpha_{22} \\ \alpha_{23} \end{Bmatrix} = \begin{Bmatrix} \dot{\omega}_{21} \\ \dot{\omega}_{22} \\ \dot{\omega}_{23} \end{Bmatrix} = [R_2] \begin{Bmatrix} C_{12} \ddot{\theta}_{11} - S_{12} \dot{\theta}_{11} \dot{\theta}_{12} \\ -S_{12} \ddot{\theta}_{11} - C_{12} \dot{\theta}_{11} \dot{\theta}_{12} \\ \ddot{\theta}_{12} \end{Bmatrix} + [\dot{R}_2] \begin{Bmatrix} C_{12} \dot{\theta}_{11} \\ -S_{12} \dot{\theta}_{11} \\ \dot{\theta}_{12} \end{Bmatrix} + \begin{Bmatrix} C_{22} \ddot{\theta}_{21} - S_{22} \dot{\theta}_{21} \dot{\theta}_{22} \\ -S_{22} \ddot{\theta}_{21} - C_{22} \dot{\theta}_{21} \dot{\theta}_{22} \\ \ddot{\theta}_{22} \end{Bmatrix}$$

$$[\dot{R}_2] = \frac{d}{dt} \begin{bmatrix} C_{22} & C_{21} S_{22} & S_{21} S_{22} \\ -S_{22} & C_{21} C_{22} & S_{21} C_{22} \\ 0 & -S_{21} & C_{21} \end{bmatrix} = \begin{bmatrix} -S_{22} \dot{\theta}_{22} & -S_{21} S_{22} \dot{\theta}_{21} + C_{21} C_{22} \dot{\theta}_{22} & C_{21} S_{22} \dot{\theta}_{21} + S_{21} C_{22} \dot{\theta}_{22} \\ -C_{22} \dot{\theta}_{22} & -S_{21} C_{22} \dot{\theta}_{21} - C_{21} S_{22} \dot{\theta}_{22} & C_{21} C_{22} \dot{\theta}_{21} - S_{21} S_{22} \dot{\theta}_{22} \\ 0 & -C_{21} \dot{\theta}_{21} & -S_{21} \dot{\theta}_{21} \end{bmatrix}$$

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