ME 6590 Multibody Dynamics Matrices and Second Order Dyadics

Dyads and Dyadics

o A dyad is a vector-vector product that has the following properties

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(\underline{a}\underline{b}) \cdot \underline{c} = (\underline{b} \cdot \underline{c})\underline{a} \qquad \text{(a vector parallel to } \underline{a})
\underline{c} \cdot (\underline{a}\underline{b}) = (\underline{c} \cdot \underline{a})\underline{b} \qquad \text{(a vector parallel to } \underline{b})
(\underline{a}\underline{b} + \underline{c}\underline{d}) \cdot \underline{e} = (\underline{b} \cdot \underline{e})\underline{a} + (\underline{d} \cdot \underline{e})\underline{c} \qquad \text{(a vector with components along } \underline{a} \text{ and } \underline{c})
\underline{e} \cdot (\underline{a}\underline{b} + \underline{c}\underline{d}) = (\underline{e} \cdot \underline{a})\underline{b} + (\underline{e} \cdot \underline{c})\underline{d} \qquad \text{(a vector with components along } \underline{b} \text{ and } \underline{d})
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- o *Dyadics* are *linear combinations* of dyads. A common example is the *inertia dyadic*.
- The *inertia dyadic* of a body about a set of *body-fixed* axes passing through its *mass-center G* may written as

$$I_{\mathcal{Z}G} = \sum_{i=1}^{3} \sum_{j=1}^{3} I_{ij}^{\prime G} \, \varrho_{i} \, \varrho_{j}$$
(1)

- Here, each of the *unit vector* products $e_i e_j$ (i, j = 1, 2, 3) are called *dyads*.
- O The *inertia values* $I_{ij}^{'G}$ form the elements of the *inertia matrix* and are called the *components* of the *dyadic* in the body-fixed reference frame $B:(\underline{e}_1,\underline{e}_2,\underline{e}_3)$.
- O Like vectors, dyadics can be represented by *different components* in *different reference frames*. Consider the dyadic \underline{A} and its representations in two different reference frames $R:(\underline{n}_1,\underline{n}_2,\underline{n}_3)$ and $S:(\underline{e}_1,\underline{e}_2,\underline{e}_3)$

$$\left[\underbrace{\underset{\approx}{A} = \sum_{k,\ell} a_{k\ell}^R \, n_k \, n_\ell}_{k\ell} = \sum_{i,j} a_{ij}^S \, \underline{e}_i \, \underline{e}_j \right] \tag{2}$$

O Here, $a_{k\ell}^R(k,\ell=1,2,3)$ represent the *components* of A_{ij} in $R:(n_1,n_2,n_3)$, and $a_{ij}^S(i,j=1,2,3)$ represent the *components* in $S:(e_1,e_2,e_3)$.

Relationship between Dyadic Components in Different Frames

- O This section formulates a *relationship between any two sets of components* of a dyadic. In the development of that relationship, it is assumed that the matrix $[R]^T$ transforms vectors and their components from frame $S:(\underline{e}_1,\underline{e}_2,\underline{e}_3)$ into frame $R:(\underline{n}_1,\underline{n}_2,\underline{n}_3)$.
- \circ The components of $\frac{A}{z}$ in two different reference frames may be related by noting

$$\sum_{i,j} a_{ij}^{S} \, \varrho_{i} \, \varrho_{j} = \sum_{i,j} a_{ij}^{S} \left(\sum_{k} R_{ik} \, n_{k} \right) \left(\sum_{\ell} R_{j\ell} \, n_{\ell} \right)$$

$$= \sum_{k,\ell} \left(\sum_{i,j} \left(a_{ij}^{S} R_{ik} R_{j\ell} \right) \right) n_{k} \, n_{\ell}$$

$$= \sum_{k,\ell} a_{k\ell}^{R} \, n_{k} \, n_{\ell}$$

o Comparing the last two equations, we note that

$$a_{k\ell}^{R} = \sum_{i,j} a_{ij}^{S} R_{ik} R_{j\ell} = \sum_{i,j} R_{ki}^{T} a_{ij}^{S} R_{j\ell}$$

Or, in matrix form

$$\overline{\left[A_{R}\right] = \left[R\right]^{T} \left[A_{S}\right] \left[R\right]} \tag{3}$$

This result can be *applied* to the *inertia matrix* of rigid bodies. Given $[I'_G]$ the inertia matrix of a body about a set of *body-fixed axes* passing through the *mass-center G*, we can calculate the inertia matrix $[I''_G]$ about any other set of axes passing through G by using Eq. (3).

$$[I_G''] = [R]^T [I_G'][R]$$

O Here $[R]^T$ represents the transformation matrix that converts vector components in the "prime" system to vector components in the "double prime" system.