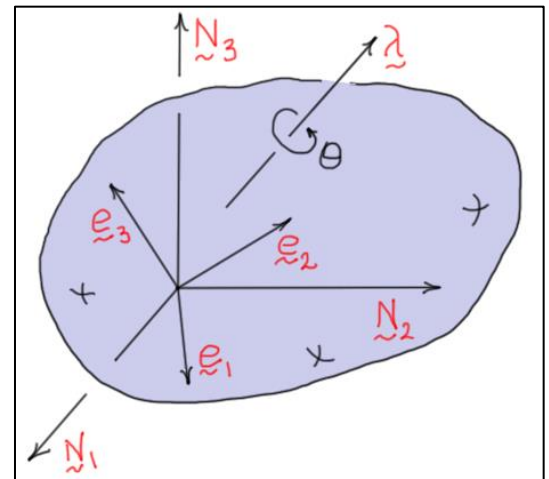


ME 5550 Intermediate Dynamics
Orientation of a Rigid Body Using Euler Parameters

Euler's Theorem on Rotation

Consider the *rigid body* shown in the figure. Let $R:(\underline{N}_1, \underline{N}_2, \underline{N}_3)$ represent the *base reference frame* and $B:(\underline{e}_1, \underline{e}_2, \underline{e}_3)$ represent the *body-fixed reference frame* and assume *initially* that the two frames are *aligned*. ***Euler's Theorem on Rotation*** states that body $B:(\underline{e}_1, \underline{e}_2, \underline{e}_3)$ can be moved into *any arbitrary orientation* relative to the base frame by a rotation about a *single axis*. In the diagram, θ represents the angle of rotation, and the *unit vector* $\underline{\lambda}$ represents the direction (or axis) of rotation.



Euler Parameters

The unit vector $\underline{\lambda}$ and the angle θ can be related to a set of *four parameters* called the ***Euler parameters***. First, let $\underline{\lambda}$ be expressed in terms of the base-frame unit vectors as

$$\underline{\lambda} = \lambda_1 \underline{N}_1 + \lambda_2 \underline{N}_2 + \lambda_3 \underline{N}_3$$

Then, the four Euler parameters are defined as follows.

$$\begin{cases} \varepsilon_1 = \lambda_1 \sin(\theta/2) \\ \varepsilon_2 = \lambda_2 \sin(\theta/2) \\ \varepsilon_3 = \lambda_3 \sin(\theta/2) \\ \varepsilon_4 = \cos(\theta/2) \end{cases} \quad \text{(four Euler parameters)}$$

Notes

1. The Euler parameters are *not independent*, because $\boxed{\varepsilon_1^2 + \varepsilon_2^2 + \varepsilon_3^2 + \varepsilon_4^2 = 1}$.
2. It can be shown that the *unit vectors* in the two reference frames can be related as follows.

$$\boxed{\begin{Bmatrix} \tilde{N}_1 \\ \tilde{N}_2 \\ \tilde{N}_3 \end{Bmatrix} = [R] \begin{Bmatrix} e_1 \\ e_2 \\ e_3 \end{Bmatrix}}$$

with

$$\boxed{[R] = \begin{bmatrix} (\varepsilon_1^2 - \varepsilon_2^2 - \varepsilon_3^2 + \varepsilon_4^2) & 2(\varepsilon_1\varepsilon_2 - \varepsilon_3\varepsilon_4) & 2(\varepsilon_1\varepsilon_3 + \varepsilon_2\varepsilon_4) \\ 2(\varepsilon_1\varepsilon_2 + \varepsilon_3\varepsilon_4) & (-\varepsilon_1^2 + \varepsilon_2^2 - \varepsilon_3^2 + \varepsilon_4^2) & 2(\varepsilon_2\varepsilon_3 - \varepsilon_1\varepsilon_4) \\ 2(\varepsilon_1\varepsilon_3 - \varepsilon_2\varepsilon_4) & 2(\varepsilon_2\varepsilon_3 + \varepsilon_1\varepsilon_4) & (-\varepsilon_1^2 - \varepsilon_2^2 + \varepsilon_3^2 + \varepsilon_4^2) \end{bmatrix}}$$

3. If the *angular velocity* of the body is resolved into *components* in the *base reference frame*, that is, ${}^R\omega_B = \omega_1\tilde{N}_1 + \omega_2\tilde{N}_2 + \omega_3\tilde{N}_3$, then it can also be shown that

$$\boxed{\begin{Bmatrix} \omega_1 \\ \omega_2 \\ \omega_3 \\ 0 \end{Bmatrix} = 2[E] \begin{Bmatrix} \dot{\varepsilon}_1 \\ \dot{\varepsilon}_2 \\ \dot{\varepsilon}_3 \\ \dot{\varepsilon}_4 \end{Bmatrix}} \quad \text{or} \quad \boxed{\begin{Bmatrix} \dot{\varepsilon}_1 \\ \dot{\varepsilon}_2 \\ \dot{\varepsilon}_3 \\ \dot{\varepsilon}_4 \end{Bmatrix} = \frac{1}{2}[E]^T \begin{Bmatrix} \omega_1 \\ \omega_2 \\ \omega_3 \\ 0 \end{Bmatrix}}$$

with

$$\boxed{[E] = \begin{bmatrix} \varepsilon_4 & -\varepsilon_3 & \varepsilon_2 & -\varepsilon_1 \\ \varepsilon_3 & \varepsilon_4 & -\varepsilon_1 & -\varepsilon_2 \\ -\varepsilon_2 & \varepsilon_1 & \varepsilon_4 & -\varepsilon_3 \\ \varepsilon_1 & \varepsilon_2 & \varepsilon_3 & \varepsilon_4 \end{bmatrix}} \quad \text{and} \quad \boxed{[E]^{-1} = [E]^T}$$

Similar expressions are true for the *angular velocity components* about the *body-fixed axes*. Note that $[E]$ is an *orthogonal matrix*.

4. Note that *no singularities* exist in the kinematic equations shown above, so many computer programs use Euler parameters to *avoid computational singularities*; however, they may *communicate* with the analyst using orientation angles which are easier to visualize.