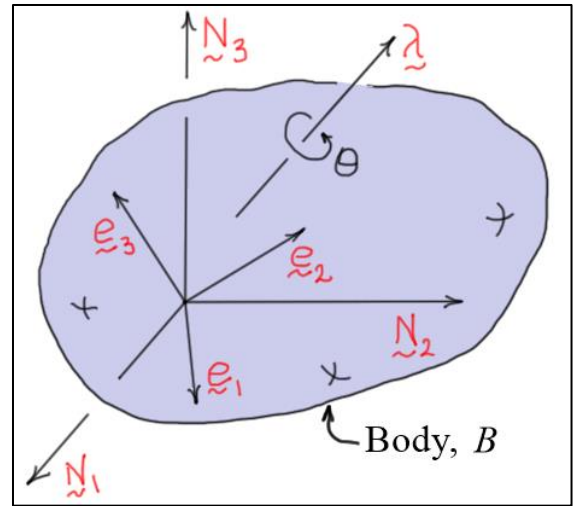


Intermediate Dynamics

Orientation of a Rigid Body Using Euler Parameters

Euler's Theorem on Rotation

Consider the *rigid body* shown in the figure. Let $R:(\underline{N}_1, \underline{N}_2, \underline{N}_3)$ represent the *base reference frame* and $B:(\underline{e}_1, \underline{e}_2, \underline{e}_3)$ represent the *body-fixed reference frame* and assume *initially* the two frames are *aligned*. **Euler's Theorem on Rotation** states that body $B:(\underline{e}_1, \underline{e}_2, \underline{e}_3)$ can be moved into *any arbitrary orientation* relative to the base frame by a *single rotation* about *some axis*.



In the diagram, θ represents the angle of rotation, and the *unit vector* $\underline{\lambda}$ represents the direction (or axis) of rotation.

Euler Parameters

The unit vector $\underline{\lambda}$ and the angle θ can be related to a set of *four parameters* called the **Euler parameters**. First, let $\underline{\lambda}$ be expressed in terms of the base-frame unit vectors as

$$\underline{\lambda} = \lambda_1 \underline{N}_1 + \lambda_2 \underline{N}_2 + \lambda_3 \underline{N}_3$$

Then, the four Euler parameters are defined as follows.

$$\begin{aligned} \varepsilon_1 &= \lambda_1 \sin(\theta/2) \\ \varepsilon_2 &= \lambda_2 \sin(\theta/2) \\ \varepsilon_3 &= \lambda_3 \sin(\theta/2) \\ \varepsilon_4 &= \cos(\theta/2) \end{aligned} \quad (\text{four Euler parameters})$$

Notes

1. The Euler parameters are *not independent*, because $\varepsilon_1^2 + \varepsilon_2^2 + \varepsilon_3^2 + \varepsilon_4^2 = 1$.

2. It can be shown that the **unit vectors** in the two reference frames can be related as follows.

$$\begin{Bmatrix} \tilde{N}_1 \\ \tilde{N}_2 \\ \tilde{N}_3 \end{Bmatrix} = [R] \begin{Bmatrix} \underline{e}_1 \\ \underline{e}_2 \\ \underline{e}_3 \end{Bmatrix}$$

with

$$[R] = \begin{bmatrix} (\varepsilon_1^2 - \varepsilon_2^2 - \varepsilon_3^2 + \varepsilon_4^2) & 2(\varepsilon_1\varepsilon_2 - \varepsilon_3\varepsilon_4) & 2(\varepsilon_1\varepsilon_3 + \varepsilon_2\varepsilon_4) \\ 2(\varepsilon_1\varepsilon_2 + \varepsilon_3\varepsilon_4) & (-\varepsilon_1^2 + \varepsilon_2^2 - \varepsilon_3^2 + \varepsilon_4^2) & 2(\varepsilon_2\varepsilon_3 - \varepsilon_1\varepsilon_4) \\ 2(\varepsilon_1\varepsilon_3 - \varepsilon_2\varepsilon_4) & 2(\varepsilon_2\varepsilon_3 + \varepsilon_1\varepsilon_4) & (-\varepsilon_1^2 - \varepsilon_2^2 + \varepsilon_3^2 + \varepsilon_4^2) \end{bmatrix}$$

3. If the **angular velocity** of the body is resolved into **components** in the **base reference frame**,

that is, ${}^R\omega_B = \omega_1\tilde{N}_1 + \omega_2\tilde{N}_2 + \omega_3\tilde{N}_3$, then it can also be shown that

$$\begin{Bmatrix} \omega_1 \\ \omega_2 \\ \omega_3 \\ 0 \end{Bmatrix} = 2[E] \begin{Bmatrix} \dot{\varepsilon}_1 \\ \dot{\varepsilon}_2 \\ \dot{\varepsilon}_3 \\ \dot{\varepsilon}_4 \end{Bmatrix} \quad \text{or} \quad \begin{Bmatrix} \dot{\varepsilon}_1 \\ \dot{\varepsilon}_2 \\ \dot{\varepsilon}_3 \\ \dot{\varepsilon}_4 \end{Bmatrix} = \frac{1}{2}[E]^T \begin{Bmatrix} \omega_1 \\ \omega_2 \\ \omega_3 \\ 0 \end{Bmatrix}$$

with

$$[E] = \begin{bmatrix} \varepsilon_4 & -\varepsilon_3 & \varepsilon_2 & -\varepsilon_1 \\ \varepsilon_3 & \varepsilon_4 & -\varepsilon_1 & -\varepsilon_2 \\ -\varepsilon_2 & \varepsilon_1 & \varepsilon_4 & -\varepsilon_3 \\ \varepsilon_1 & \varepsilon_2 & \varepsilon_3 & \varepsilon_4 \end{bmatrix} \quad \text{and} \quad [E]^{-1} = [E]^T$$

Similar expressions are true for the **angular velocity components** about the **body-fixed axes**. Note that $[E]$ is an **orthogonal matrix**.

4. Note that **no singularities** exist in the kinematic equations shown above, so many computer programs use Euler parameters (or Euler-like parameters) to **avoid computational singularities**. They may, however, **communicate** with the analyst using orientation angles which are easier to visualize.