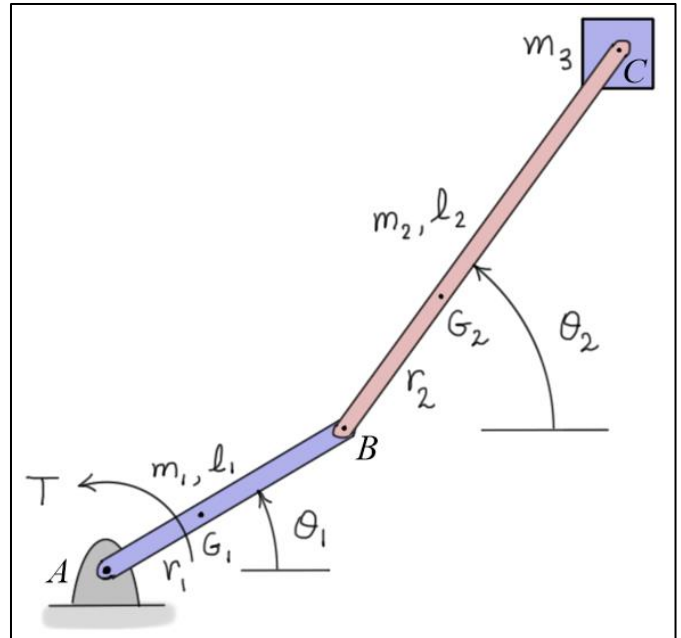


## Intermediate Dynamics

### Equations of Motion of a Slider-Crank Mechanism

The equations of motion of a *slider-crank mechanism* can be formulated in various ways. Here, the equations of motion are formulated for the system shown in the diagram. The necessary constraints to form the slider-crank mechanism are imposed using Lagrange's equations with Lagrange multipliers. The system shown consists of *two links* and an *end mass*. A torque  $T$  drives link *AB* on the shaft at *A*. It is assumed that the *end mass*  $m_3$  *translates but does not rotate*.



#### Constraint

The system shown can be converted into a simple *slider-crank mechanism* with *zero offset*, by imposing the following configuration constraint.

$$\boxed{l_1 \sin(\theta_1) + l_2 \sin(\theta_2) = 0} \quad (1)$$

This constraint can be put into *standard form* by *differentiating* it with respect to *time* to get

$$\boxed{(\ell_1 \cos(\theta_1))\dot{\theta}_1 + (\ell_2 \cos(\theta_2))\dot{\theta}_2 = 0} \quad (2)$$

#### Equations of Motion

The equations of motion of the slider-crank mechanism can be developed using *Lagrange's equations* with a single *Lagrange multiplier*. That is,

$$\boxed{\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\theta}_i} \right) - \frac{\partial L}{\partial \theta_i} = F_{\theta_i} + \lambda_1 a_{1i}} \quad (i = 1, 2) \quad (3)$$

The coefficients  $a_{1i}$  ( $i = 1, 2$ ) are found by comparing Eq. (2) with the standard constraint equation form.

$$\boxed{a_{11} = \ell_1 \cos(\theta_1)} \quad \boxed{a_{12} = \ell_2 \cos(\theta_2)} \quad (4)$$

Including the *kinetic energies* of the *three bodies* and the *potential energies* associated with the *weight forces*, it can be shown that

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\theta}_1} \right) - \frac{\partial L}{\partial \theta_1} = (m_1 r_1^2 + I_1 + m_2 \ell_1^2 + m_3 \ell_1^2) \ddot{\theta}_1 + (m_2 \ell_1 r_2 + m_3 \ell_1 \ell_2) \cos(\theta_2 - \theta_1) \ddot{\theta}_2 - (m_2 \ell_1 r_2 + m_3 \ell_1 \ell_2) \sin(\theta_2 - \theta_1) \dot{\theta}_2^2 + (m_1 r_1 + m_2 \ell_1 + m_3 \ell_1) g \cos(\theta_1) \quad (5)$$

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\theta}_2} \right) - \frac{\partial L}{\partial \theta_2} = (m_2 \ell_1 r_2 + m_3 \ell_1 \ell_2) \cos(\theta_2 - \theta_1) \ddot{\theta}_1 + (m_2 r_2^2 + I_2 + m_3 \ell_2^2) \ddot{\theta}_2 + (m_2 \ell_1 r_2 + m_3 \ell_1 \ell_2) \sin(\theta_2 - \theta_1) \dot{\theta}_1^2 + (m_2 r_2 + m_3 \ell_2) g \cos(\theta_2) \quad (6)$$

The *contribution* of the *driving torque* to the equations of motion are

$$F_{\theta_1} = T \tilde{k} \cdot \left[ \frac{\partial \varpi_1}{\partial \dot{\theta}_1} \right] = T \quad F_{\theta_2} = T \tilde{k} \cdot \left[ \frac{\partial \varpi_1}{\partial \dot{\theta}_2} \right] = 0 \quad (7)$$

Using these results in the equations of motion (Eq. (3)) gives the equations of motion.

$$(m_1 r_1^2 + I_1 + m_2 \ell_1^2 + m_3 \ell_1^2) \ddot{\theta}_1 + (m_2 \ell_1 r_2 + m_3 \ell_1 \ell_2) \cos(\theta_2 - \theta_1) \ddot{\theta}_2 - (m_2 \ell_1 r_2 + m_3 \ell_1 \ell_2) \sin(\theta_2 - \theta_1) \dot{\theta}_2^2 + (m_1 r_1 + m_2 \ell_1 + m_3 \ell_1) g \cos(\theta_1) = T + \lambda_1 \ell_1 \cos(\theta_1) \quad (8)$$

$$(m_2 \ell_1 r_2 + m_3 \ell_1 \ell_2) \cos(\theta_2 - \theta_1) \ddot{\theta}_1 + (m_2 r_2^2 + I_2 + m_3 \ell_2^2) \ddot{\theta}_2 + (m_2 \ell_1 r_2 + m_3 \ell_1 \ell_2) \sin(\theta_2 - \theta_1) \dot{\theta}_1^2 + (m_2 r_2 + m_3 \ell_2) g \cos(\theta_2) = \lambda_1 \ell_2 \cos(\theta_2) \quad (9)$$

Eqs. (8) and (9) must be solved along with the constraint equation to find  $\theta_1$ ,  $\theta_2$ , and  $\lambda_1$ .

Differentiating Eq. (2) gives

$$(\ell_1 \cos(\theta_1)) \ddot{\theta}_1 + (\ell_2 \cos(\theta_2)) \ddot{\theta}_2 - (\ell_1 \sin(\theta_1)) \dot{\theta}_1^2 - (\ell_2 \sin(\theta_2)) \dot{\theta}_2^2 = 0 \quad (10)$$

Eqs. (8), (9), and (10) form a set of *three coupled, second-order, differential/algebraic equations* that can be solved for the *three unknowns*  $\theta_1$ ,  $\theta_2$ , and  $\lambda_1$  as functions of time given the *driving torque*  $T$  and an *initial position*.

The approach taken here represents just *one way* of breaking the system down and then putting it back together with configuration constraints. Question: What other systems could be joined to form the slider-crank mechanism?