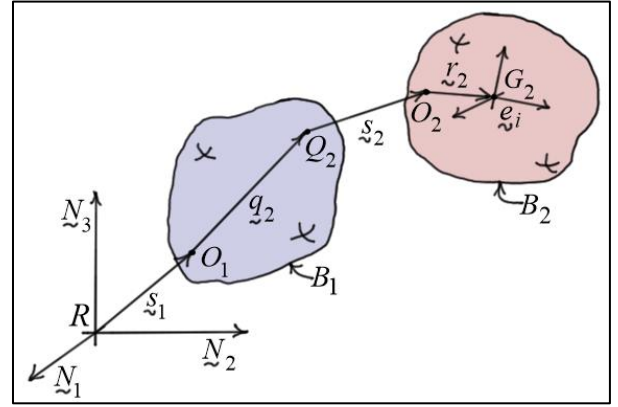


Multibody Dynamics

Accelerations Using Relative Coordinates

Velocities

Consider again the *two-body system* shown. In previous notes, it was found that the *fixed-frame components* of ${}^R v_{G_2}$ the *velocity* of G_2 in the fixed frame R can be written as



$$\{v_{G_2}\} = \{\dot{s}_1\} - \left([\tilde{q}_2] + [\tilde{s}_2] + [\tilde{r}_2] \right) \{\omega_{B_1}\} + [R_{B_1}]^T \{\dot{s}'_2\} - [\tilde{r}_2] [R_{B_1}]^T \{\hat{\omega}_{B_2}\} \quad (1)$$

or

$$\begin{aligned} \{v_{G_2}\} = & \{\dot{s}_1\} - [R_{B_1}]^T \left([\tilde{q}'_2] + [\tilde{s}'_2] \right) \{\omega'_{B_1}\} - [R_{B_2}]^T [\tilde{r}'_2] [{}^{B_2}R_{B_1}]^T \{\omega'_{B_1}\} \\ & + [R_{B_1}]^T \{\dot{s}'_2\} - [R_{B_2}]^T [\tilde{r}'_2] \{\hat{\omega}'_{B_2}\} \end{aligned} \quad (2)$$

The first equation is written in terms of the *fixed-frame components* of ${}^R \omega_{B_1}$ and the B_1 *frame components* of ${}^{B_1} \omega_{B_2}$, and the second is written in terms of the *body-frame components* of the angular velocities.

Accelerations

Case 1: Fixed-Frame Angular Velocity Components as Generalized Speeds

If the generalized speeds are defined to be the elements of $\{\dot{s}_1\}$, $\{\dot{s}'_2\}$, $\{\omega_{B_1}\}$, and $\{\hat{\omega}_{B_2}\}$, then write

$$\boxed{\{v_{G_2}\} = [v_{G_2, \dot{s}_1}] \{\dot{s}_1\} + [v_{G_2, \omega_{B_1}}] \{\omega_{B_1}\} + [v_{G_2, \dot{s}'_2}] \{\dot{s}'_2\} + [v_{G_2, \hat{\omega}_{B_2}}] \{\hat{\omega}_{B_2}\}} \quad (3)$$

Here,

$[v_{G_2, \dot{s}_1}]$ is the 3×3 identity matrix

$$[v_{G_2, \omega_{B_1}}] = - \left([\tilde{q}_2] + [\tilde{s}_2] + [\tilde{r}_2] \right)$$

$$[v_{G_2, \dot{s}'_2}] = [R_{B_1}]^T$$

$$[v_{G_2, \hat{\omega}_{B_2}}] = - [\tilde{r}_2] [R_{B_1}]^T$$

Differentiating Eq. (3), the *fixed-frame components* of the *acceleration* of G_2 can then be written as

$$\boxed{\begin{aligned} \{a_{G_2}\} &= [v_{G_2, \dot{s}_1}] \{\ddot{s}_1\} + [v_{G_2, \omega_{B_1}}] \{\dot{\omega}_{B_1}\} + [\dot{v}_{G_2, \omega_{B_1}}] \{\omega_{B_1}\} \\ &+ [v_{G_2, \dot{s}'_2}] \{\ddot{s}'_2\} + [\dot{v}_{G_2, \dot{s}'_2}] \{\dot{s}'_2\} + [v_{G_2, \hat{\omega}_{B_2}}] \{\dot{\hat{\omega}}_{B_2}\} + [\dot{v}_{G_2, \hat{\omega}_{B_2}}] \{\hat{\omega}_{B_2}\} \end{aligned}} \quad (4)$$

where the time derivatives of the partial velocity matrices can be written as follows.

$$a) \quad [\dot{v}_{G_2, \omega_{B_1}}] = -\left([\ddot{q}_2] + [\dot{s}_2] + [\dot{r}_2]\right)$$

The components of the vectors $\{\dot{q}_2\}$, $\{\dot{s}_2\}$, and $\{\dot{r}_2\}$ are elements of the matrices $[\ddot{q}_2]$, $[\dot{s}_2]$, and $[\dot{r}_2]$. These components are found as follows.

$$\{\dot{q}_2\} = \frac{d}{dt} \left([R_{B_1}]^T \{q'_2\} \right) = [\dot{R}_{B_1}]^T \{q'_2\} = [\tilde{\omega}_{B_1}] [R_{B_1}]^T \{q'_2\} \quad (5)$$

$$\begin{aligned} \{\dot{s}_2\} &= \frac{d}{dt} \left([R_{B_1}]^T \{s'_2\} \right) = [\dot{R}_{B_1}]^T \{s'_2\} + [R_{B_1}]^T \{\dot{s}'_2\} \\ &= [\tilde{\omega}_{B_1}] [R_{B_1}]^T \{s'_2\} + [R_{B_1}]^T \{\dot{s}'_2\} \end{aligned} \quad (6)$$

$$\{\dot{r}_2\} = \frac{d}{dt} \left([R_{B_2}]^T \{r'_2\} \right) = [\dot{R}_{B_2}]^T \{r'_2\} = [\tilde{\omega}_{B_2}] [R_{B_2}]^T \{r'_2\} \quad (7)$$

$$b) \quad [\dot{v}_{G_2, \dot{s}'_2}] = [\dot{R}_{B_1}]^T = [\tilde{\omega}_{B_1}] [R_{B_1}]^T \quad (8)$$

$$c) \quad [\dot{v}_{G_2, \hat{\omega}_{B_2}}] = -[\dot{r}_2] [R_{B_1}]^T - [\dot{r}_2] [\dot{R}_{B_1}]^T = -[\dot{r}_2] [R_{B_1}]^T - [\dot{r}_2] [\tilde{\omega}_{B_1}] [R_{B_1}]^T \quad (9)$$

The elements of $[\dot{r}_2]$ are calculated in Eq. (7) above.

Case 2: Body-Frame Angular Velocity Components as Generalized Speeds

If the *generalized speeds* are defined to be the components of $\{\dot{s}_1\}$, $\{\dot{s}'_2\}$, $\{\omega'_{B_1}\}$, and $\{\hat{\omega}'_{B_2}\}$,

then

$$\boxed{\{v_{G_2}\} = [v_{G_2, \dot{s}_1}] \{\dot{s}_1\} + [v_{G_2, \omega'_{B_1}}] \{\omega'_{B_1}\} + [v_{G_2, \dot{s}'_2}] \{\dot{s}'_2\} + [v_{G_2, \hat{\omega}'_{B_2}}] \{\hat{\omega}'_{B_2}\}} \quad (10)$$

Here,

$\begin{bmatrix} v_{G_2, \dot{s}_1} \end{bmatrix}$ is the 3×3 identity matrix

$$\begin{bmatrix} v_{G_2, \omega'_{B_1}} \end{bmatrix} = -\begin{bmatrix} R_{B_1} \end{bmatrix}^T \left(\begin{bmatrix} \tilde{q}'_2 \end{bmatrix} + \begin{bmatrix} \tilde{s}'_2 \end{bmatrix} \right) - \begin{bmatrix} R_{B_2} \end{bmatrix}^T \begin{bmatrix} \tilde{r}'_2 \end{bmatrix} \begin{bmatrix} {}^{B_2}R_{B_1} \end{bmatrix}^T$$

$$\begin{bmatrix} v_{G_2, \dot{s}'_2} \end{bmatrix} = \begin{bmatrix} R_{B_1} \end{bmatrix}^T$$

$$\begin{bmatrix} v_{G_2, \hat{\omega}'_{B_2}} \end{bmatrix} = -\begin{bmatrix} R_{B_2} \end{bmatrix}^T \begin{bmatrix} \tilde{r}'_2 \end{bmatrix}$$

Differentiating Eq. (10), the *fixed-frame components* of the *acceleration* of G_2 can then be written as

$$\boxed{\begin{aligned} \{a_{G_2}\} &= \begin{bmatrix} v_{G_2, \dot{s}_1} \end{bmatrix} \{\ddot{s}_1\} + \begin{bmatrix} v_{G_2, \omega'_{B_1}} \end{bmatrix} \{\dot{\omega}'_{B_1}\} + \begin{bmatrix} \dot{v}_{G_2, \omega'_{B_1}} \end{bmatrix} \{\omega'_{B_1}\} \\ &+ \begin{bmatrix} v_{G_2, \dot{s}'_2} \end{bmatrix} \{\ddot{s}'_2\} + \begin{bmatrix} \dot{v}_{G_2, \dot{s}'_2} \end{bmatrix} \{\dot{s}'_2\} + \begin{bmatrix} v_{G_2, \hat{\omega}'_{B_2}} \end{bmatrix} \{\dot{\hat{\omega}}'_{B_2}\} + \begin{bmatrix} \dot{v}_{G_2, \hat{\omega}'_{B_2}} \end{bmatrix} \{\hat{\omega}'_{B_2}\} \end{aligned}}$$

Here, the time derivatives of the partial velocity matrices can be written as follows.

$$\begin{aligned} \text{a) } \begin{bmatrix} \dot{v}_{G_2, \omega'_{B_1}} \end{bmatrix} &= -\begin{bmatrix} R_{B_1} \end{bmatrix}^T \begin{bmatrix} \dot{\tilde{s}}'_2 \end{bmatrix} - \begin{bmatrix} \dot{R}_{B_1} \end{bmatrix}^T \left(\begin{bmatrix} \tilde{q}'_2 \end{bmatrix} + \begin{bmatrix} \tilde{s}'_2 \end{bmatrix} \right) - \begin{bmatrix} \dot{R}_{B_2} \end{bmatrix}^T \begin{bmatrix} \tilde{r}'_2 \end{bmatrix} \begin{bmatrix} {}^{B_2}R_{B_1} \end{bmatrix}^T \\ &- \begin{bmatrix} R_{B_2} \end{bmatrix}^T \begin{bmatrix} \tilde{r}'_2 \end{bmatrix} \begin{bmatrix} {}^{B_1}\dot{R}_{B_2} \end{bmatrix} \\ &= -\begin{bmatrix} R_{B_1} \end{bmatrix}^T \begin{bmatrix} \dot{\tilde{s}}'_2 \end{bmatrix} - \begin{bmatrix} R_{B_1} \end{bmatrix}^T \begin{bmatrix} \tilde{\omega}'_{B_1} \end{bmatrix} \left(\begin{bmatrix} \tilde{q}'_2 \end{bmatrix} + \begin{bmatrix} \tilde{s}'_2 \end{bmatrix} \right) \\ &- \begin{bmatrix} R_{B_2} \end{bmatrix}^T \begin{bmatrix} \tilde{\omega}'_{B_2} \end{bmatrix} \begin{bmatrix} \tilde{r}'_2 \end{bmatrix} \begin{bmatrix} {}^{B_2}R_{B_1} \end{bmatrix}^T - \begin{bmatrix} R_{B_2} \end{bmatrix}^T \begin{bmatrix} \tilde{r}'_2 \end{bmatrix} \begin{bmatrix} \tilde{\omega}'_{B_2} \end{bmatrix} \begin{bmatrix} {}^{B_1}R_{B_2} \end{bmatrix} \end{aligned}$$

$$\text{b) } \begin{bmatrix} \dot{v}_{G_2, \dot{s}'_2} \end{bmatrix} = \begin{bmatrix} \dot{R}_{B_1} \end{bmatrix}^T = \begin{bmatrix} R_{B_1} \end{bmatrix}^T \begin{bmatrix} \tilde{\omega}'_{B_1} \end{bmatrix}$$

$$\text{c) } \begin{bmatrix} \dot{v}_{G_2, \hat{\omega}'_{B_2}} \end{bmatrix} = -\begin{bmatrix} \dot{R}_{B_2} \end{bmatrix}^T \begin{bmatrix} \tilde{r}'_2 \end{bmatrix} = -\begin{bmatrix} R_{B_2} \end{bmatrix}^T \begin{bmatrix} \tilde{\omega}'_{B_2} \end{bmatrix} \begin{bmatrix} \tilde{r}'_2 \end{bmatrix}$$

Case 3: Orientation Angle Derivatives as Generalized Speeds

If, instead, the generalized speeds are defined to be the components of $\{\dot{s}_1\}$, $\{\dot{s}_2\}$, and a set of *orientation angle derivatives*, then $\{\omega_{B_1}\}$ and $\{\hat{\omega}_{B_2}\}$ (or $\{\omega'_{B_1}\}$ and $\{\hat{\omega}'_{B_2}\}$) can be written in terms of the orientation angles. Returning to Eq. (1), for example,

$$\begin{aligned}
\{v_{G_2}\} &= \{\dot{s}_1\} - \left([\tilde{q}_2] + [\tilde{s}_2] + [\tilde{r}_2] \right) \{\omega_{B_1}\} + [R_{B_1}]^T \{\dot{s}'_2\} - [\tilde{r}_2] [R_{B_1}]^T \{\hat{\omega}_{B_2}\} \\
&= \{\dot{s}_1\} - \left([\tilde{q}_2] + [\tilde{s}_2] + [\tilde{r}_2] \right) [\omega_{B_1, \dot{\theta}_{B_1}}] \{\dot{\theta}_{B_1}\} + [R_{B_1}]^T \{\dot{s}'_2\} - [\tilde{r}_2] [\omega_{B_2, \dot{\theta}_{B_2}}] \{\dot{\theta}_{B_2}\} \quad (11) \\
&= [v_{G_2, \dot{s}_1}] \{\dot{s}_1\} + [v_{G_2, \dot{\theta}_{B_1}}] \{\dot{\theta}_{B_1}\} + [v_{G_2, \dot{s}'_2}] \{\dot{s}'_2\} + [v_{G_2, \dot{\theta}_{B_2}}] \{\dot{\theta}_{B_2}\}
\end{aligned}$$

Here,

$$\begin{aligned}
[v_{G_2, \dot{s}_1}] &\text{ is the } 3 \times 3 \text{ identity matrix} \\
[v_{G_2, \dot{\theta}_{B_1}}] &= - \left([\tilde{q}_2] + [\tilde{s}_2] + [\tilde{r}_2] \right) [\omega_{B_1, \dot{\theta}_{B_1}}] \\
[v_{G_2, \dot{s}'_2}] &= [R_{B_1}]^T \\
[v_{G_2, \dot{\theta}_{B_2}}] &= - [\tilde{r}_2] [\omega_{B_2, \dot{\theta}_{B_2}}]
\end{aligned}$$

The partial angular velocity matrices are defined in previous notes for a 1-2-3 rotation sequence.

The fixed-frame components of the acceleration of G_2 can then be written as

$$\begin{aligned}
\{a_{G_2}\} &= [v_{G_2, \dot{s}_1}] \{\ddot{s}_1\} + [v_{G_2, \dot{\theta}_{B_1}}] \{\ddot{\theta}_{B_1}\} + [\dot{v}_{G_2, \dot{\theta}_{B_1}}] \{\dot{\theta}_{B_1}\} + [v_{G_2, \dot{s}'_2}] \{\ddot{s}'_2\} + [\dot{v}_{G_2, \dot{s}'_2}] \{\dot{s}'_2\} \\
&\quad + [v_{G_2, \dot{\theta}_{B_2}}] \{\ddot{\theta}_{B_2}\} + [\dot{v}_{G_2, \dot{\theta}_{B_2}}] \{\dot{\theta}_{B_2}\} \quad (12)
\end{aligned}$$

Here,

$$\begin{aligned}
[\dot{v}_{G_2, \dot{\theta}_{B_1}}] &= - \left([\dot{\tilde{q}}_2] + [\dot{\tilde{s}}_2] + [\dot{\tilde{r}}_2] \right) [\omega_{B_1, \dot{\theta}_{B_1}}] - \left([\tilde{q}_2] + [\tilde{s}_2] + [\tilde{r}_2] \right) [\dot{\omega}_{B_1, \dot{\theta}_{B_1}}] \\
[\dot{v}_{G_2, \dot{s}'_2}] &= [\dot{R}_{B_1}]^T = [\tilde{\omega}_{B_1}] [R_{B_1}]^T \\
[\dot{v}_{G_2, \dot{\theta}_{B_2}}] &= - [\dot{\tilde{r}}_2] [\omega_{B_2, \dot{\theta}_{B_2}}] - [\tilde{r}_2] [\dot{\omega}_{B_2, \dot{\theta}_{B_2}}]
\end{aligned}$$