

# Multibody Dynamics

## Matrices and Second Order Dyadics

### Dyads and Dyadics

- A **dyad** is a **vector-vector product** that has the following properties

$$\begin{aligned}
 (\underline{a}\underline{b}) \cdot \underline{c} &= (\underline{b} \cdot \underline{c})\underline{a} && \text{(a vector parallel to } \underline{a} \text{)} \\
 \underline{c} \cdot (\underline{a}\underline{b}) &= (\underline{c} \cdot \underline{a})\underline{b} && \text{(a vector parallel to } \underline{b} \text{)} \\
 (\underline{a}\underline{b} + \underline{c}\underline{d}) \cdot \underline{e} &= (\underline{b} \cdot \underline{e})\underline{a} + (\underline{d} \cdot \underline{e})\underline{c} && \text{(a vector with components along } \underline{a} \text{ and } \underline{c} \text{)} \\
 \underline{e} \cdot (\underline{a}\underline{b} + \underline{c}\underline{d}) &= (\underline{e} \cdot \underline{a})\underline{b} + (\underline{e} \cdot \underline{c})\underline{d} && \text{(a vector with components along } \underline{b} \text{ and } \underline{d} \text{)}
 \end{aligned}$$

- **Dyadics** are **linear combinations** of dyads. A common example is the **inertia dyadic**.
- The **inertia dyadic** of a body about a set of **body-fixed** axes passing through its **mass-center**  $G$  can written as

$$\boxed{\underline{I}_G = \sum_{i=1}^3 \sum_{j=1}^3 I_{ij}^G \underline{e}_i \underline{e}_j} \tag{1}$$

- Here, each of the **unit vector** products  $\underline{e}_i \underline{e}_j$  ( $i, j = 1, 2, 3$ ) are called **dyads**.
- The **inertia values**  $I_{ij}^G$  form the elements of the **inertia matrix** and are called the **components** of the **dyadic** in the body-fixed reference frame  $B: (\underline{e}_1, \underline{e}_2, \underline{e}_3)$ .
- Like vectors, dyadics can be represented by **different components** in **different reference frames**. Consider the dyadic  $\underline{A}$  and its representations in two different reference frames  $R: (\underline{n}_1, \underline{n}_2, \underline{n}_3)$  and  $S: (\underline{e}_1, \underline{e}_2, \underline{e}_3)$

$$\boxed{\underline{A} = \sum_{k,\ell} a_{k\ell}^R \underline{n}_k \underline{n}_\ell = \sum_{i,j} a_{ij}^S \underline{e}_i \underline{e}_j} \tag{2}$$

- Here,  $a_{k\ell}^R$  ( $k, \ell = 1, 2, 3$ ) represent the **components** of  $\underline{A}$  in  $R: (\underline{n}_1, \underline{n}_2, \underline{n}_3)$ , and  $a_{ij}^S$  ( $i, j = 1, 2, 3$ ) represent the **components** in  $S: (\underline{e}_1, \underline{e}_2, \underline{e}_3)$ .

## Relationship between Dyadic Components in Different Frames

- This section formulates a **relationship between any two sets of components** of a dyadic. In the development of that relationship, it is assumed that the matrix  $[R]^T$  transforms vectors and their components from frame  $S:(e_1, e_2, e_3)$  into frame  $R:(n_1, n_2, n_3)$ .
- The components of  $A$  in **two different reference frames** can be **related** by noting

$$\begin{aligned} \sum_{i,j} a_{ij}^S e_i e_j &= \sum_{i,j} a_{ij}^S \left( \sum_k R_{ik} n_k \right) \left( \sum_\ell R_{j\ell} n_\ell \right) \\ &= \sum_{k,\ell} \left( \sum_{i,j} (a_{ij}^S R_{ik} R_{j\ell}) \right) n_k n_\ell \\ &= \sum_{k,\ell} a_{k\ell}^R n_k n_\ell \end{aligned}$$

- **Comparing the last two equations**, we note that

$$\boxed{a_{k\ell}^R = \sum_{i,j} a_{ij}^S R_{ik} R_{j\ell} = \sum_{i,j} R_{ki}^T a_{ij}^S R_{j\ell}}$$

Or, in matrix form

$$\boxed{[A_R] = [R]^T [A_S] [R]} \quad (3)$$

- This result can be **applied** to the **inertia matrix** of rigid bodies. Given  $[I'_G]$  the inertia matrix of a body about a set of **body-fixed axes** passing through the **mass-center**  $G$ , we can calculate the inertia matrix  $[I''_G]$  about any other set of axes passing through  $G$  by using Eq. (3).

$$\boxed{[I''_G] = [R]^T [I'_G] [R]}$$

- Here  $[R]^T$  represents the transformation matrix that converts vector components in the “**prime**” system to vector components in the “**double prime**” system.